

# The cuprate phase diagram: insights from neutron scattering and electrical transport

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# Acknowledgements

**Martin Greven, University of Minnesota**

## **CRYSTAL GROWTH**

C. Dorow, M. Veit, Y. Ge (UMN); X. Zhao (Jilin U., China)

## **ELECTRICAL TRANSPORT**

**N. Barišić** (CEA-Saclay, France, Vienna University of Technology, Austria)

W. Tabis, M. Veit, Y. Li (UMN)

C. Proust (Toulouse, France)

## **NEUTRON SCATTERING**

Y. Tang, C. Dorow, M. Veit (UMN)

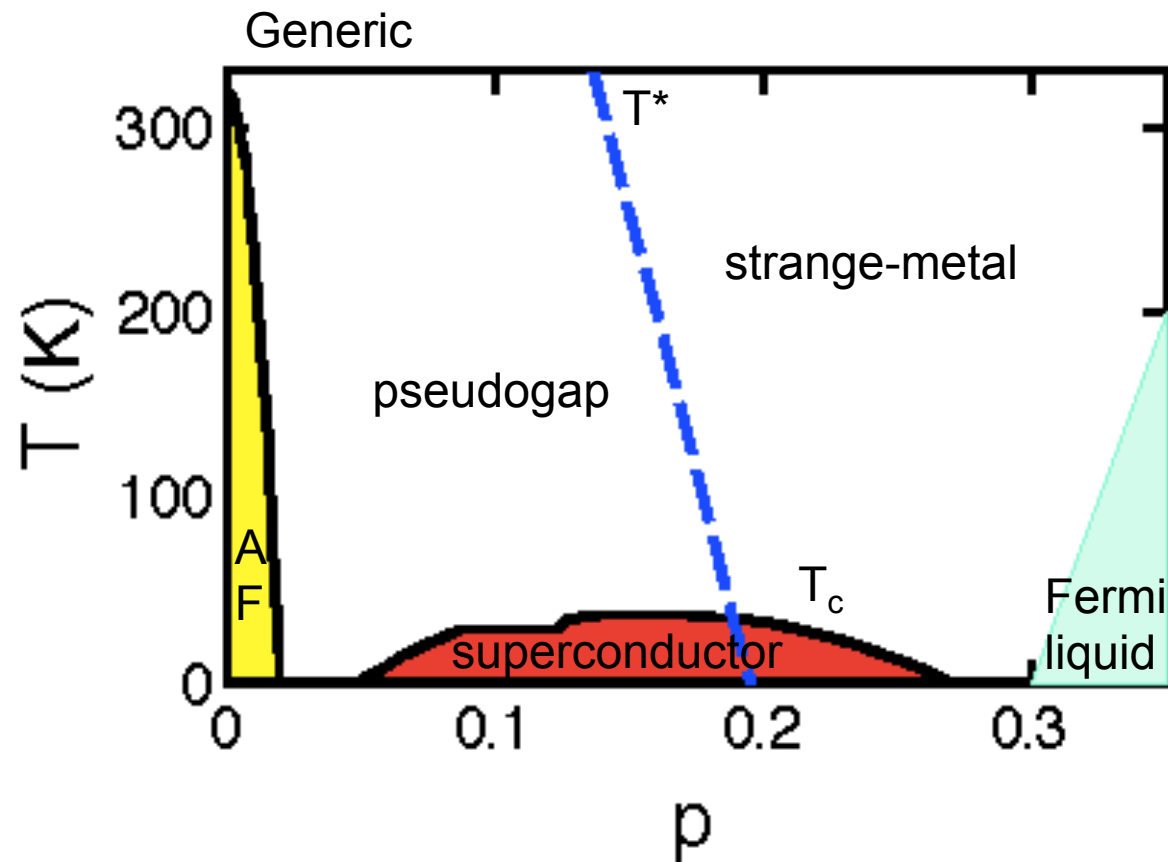
L. Mangin-Thro, Y. Sidis, P. Bourges (LLB, CEA-Saclay, France)

D. L. Abernathy, A. D. Christianson (ORNL, USA)

P. Steffens (ILL, Grenoble, France)

J.T. Park (FRM-II, Garching, Germany)

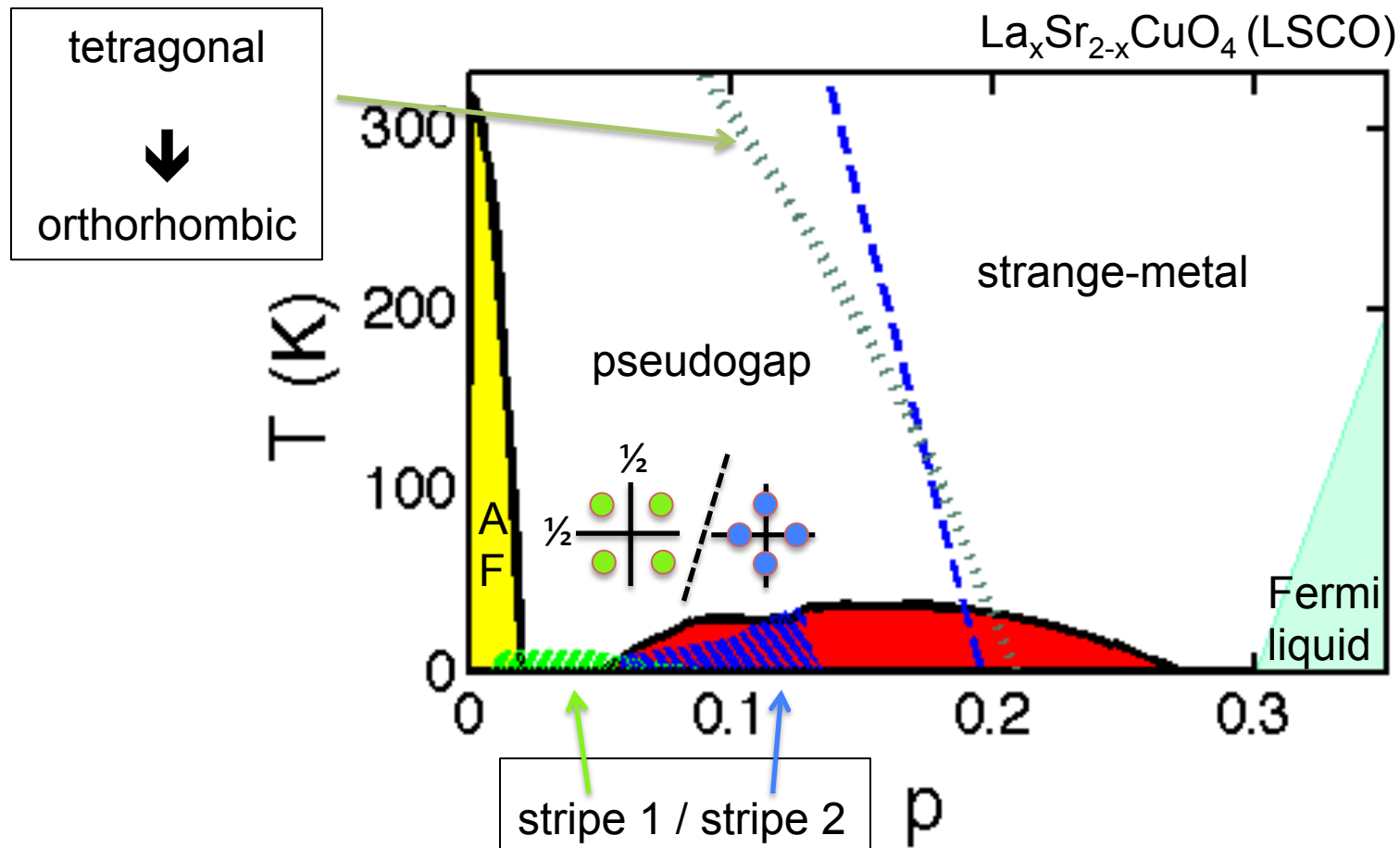
# Cuprate phase diagram



- Insulating parent compound: localized spins with antiferromagnetic order
- Fermi-liquid at very high doping

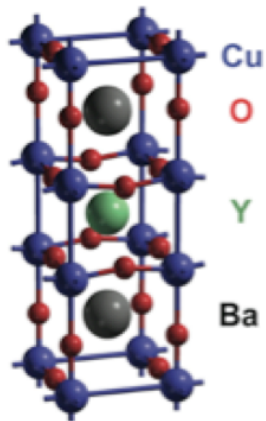
# Structural transitions and competing phases

- Why is high- $T_c$  a hard problem (experimentally)?
  - Disorder (e.g. Sr substitution) and structural transitions
  - Other ordering tendencies: **Universal?** **Integral?**



# The most studied cuprates

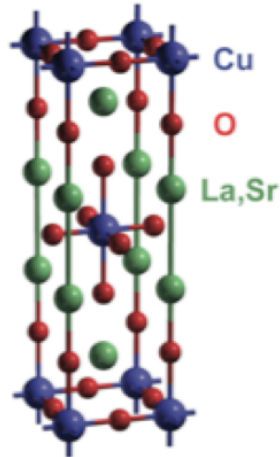
$\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$   
(YBCO)



courtesy Guichuan Yu

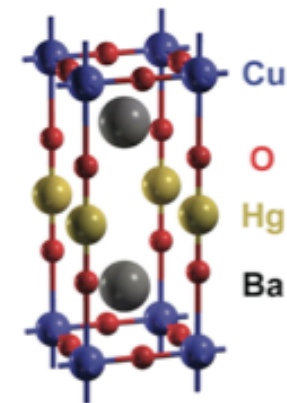
- CuO chains
- 2 x  $\text{CuO}_2$  planes
- orthorhombic
- $T_c^{\text{max}} = 93 \text{ K}$

$\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$   
(LSCO)



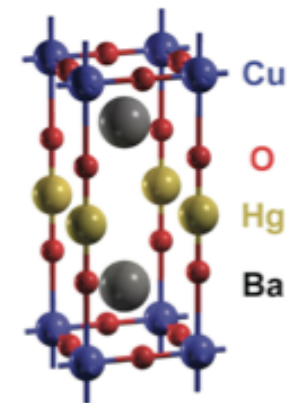
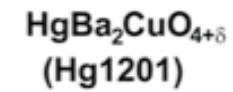
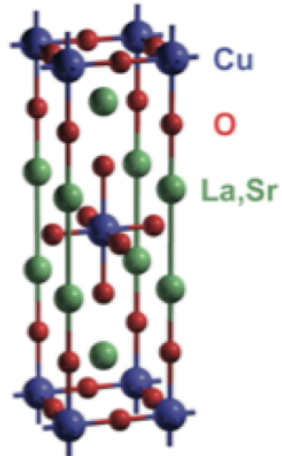
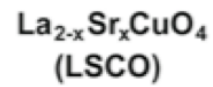
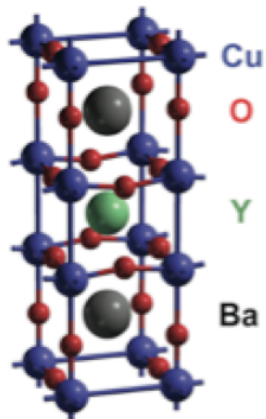
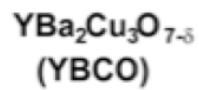
- No chains
- 1 x  $\text{CuO}_2$  plane
- orthorhombic
- $T_c^{\text{max}} = 40 \text{ K}$

$\text{HgBa}_2\text{CuO}_{4+\delta}$   
(Hg1201)

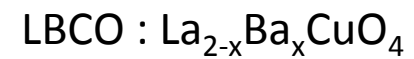
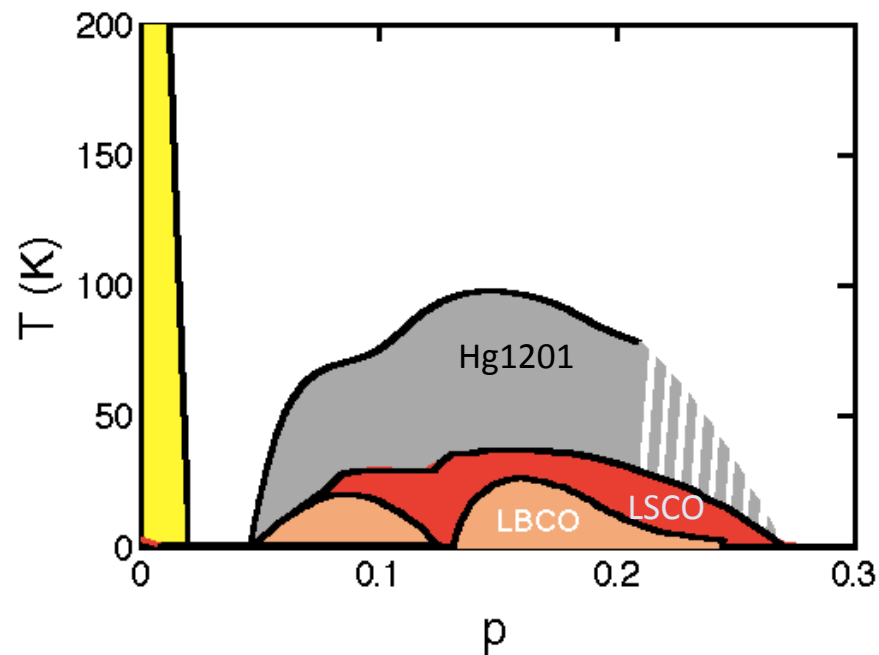


- No  $\text{CuO}_2$  chains
- 1 x CuO layer
- tetragonal
- $T_c^{\text{max}} = 98 \text{ K}$  (the highest of the single-layered cuprates)

# The most studied cuprates



courtesy Guichuan Yu

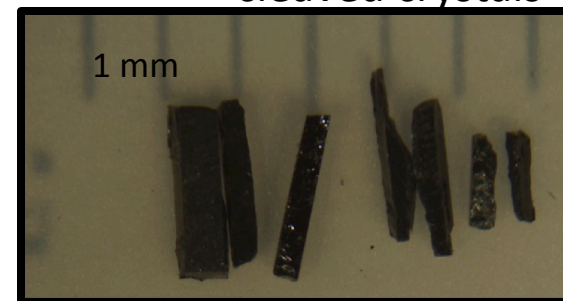


# Part 1: electrical transport



Mike Veit (senior, UMN)

cleaved crystals



*LNCMI-Toulouse, France*  
Baptiste Vignolle, Cyril Proust



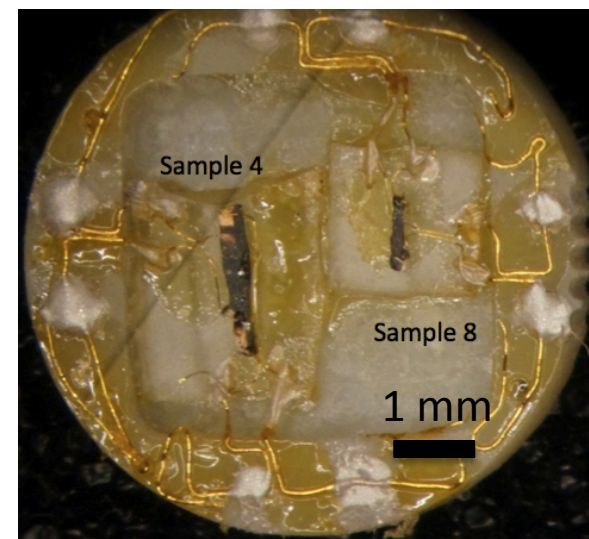
Au pads+Ag paste



*Vienna University of Technology,  
Austria*  
N. Barišić



contacted and ready for pulse field

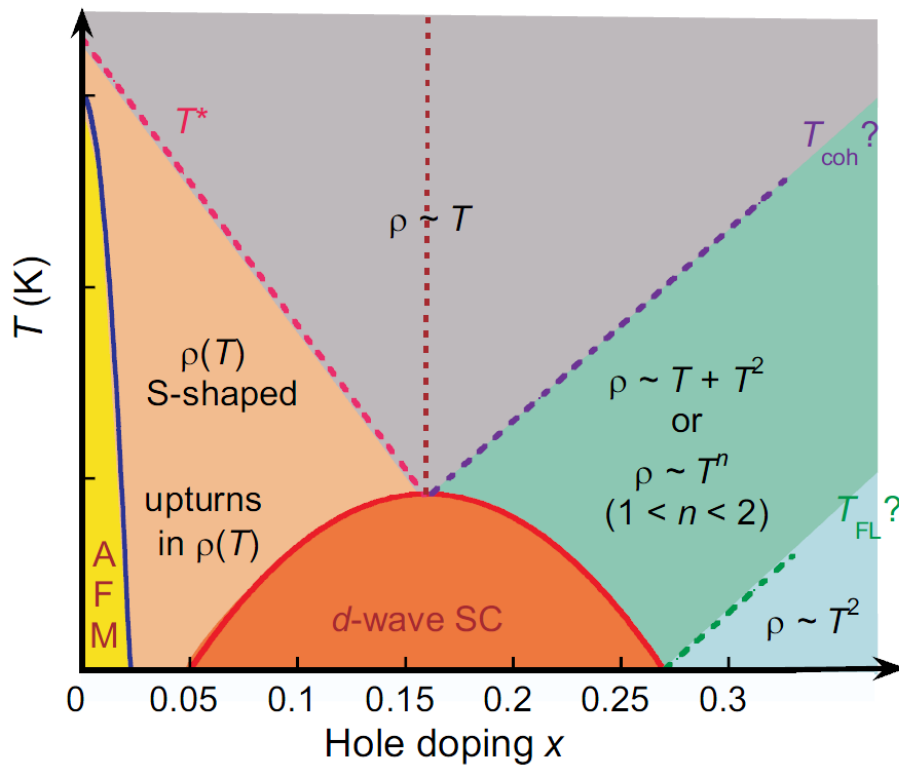


*National High Magnetic Field Lab  
Tallahassee and Los Alamos*  
A. Sheckter  
R. McDonald

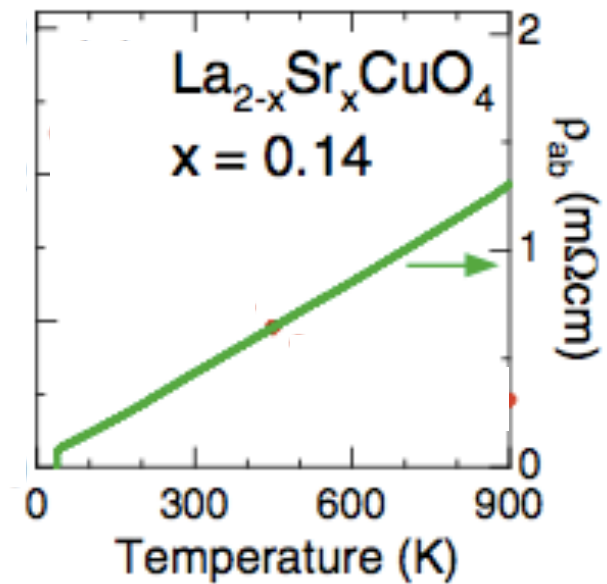


# Resistivity (strange-metal)

- T- linear resistivity to very high temperatures



N. J. Hussey Phys.: Condens. Matter. (2008)



Y. Ando *et al.*, PRL (2004)



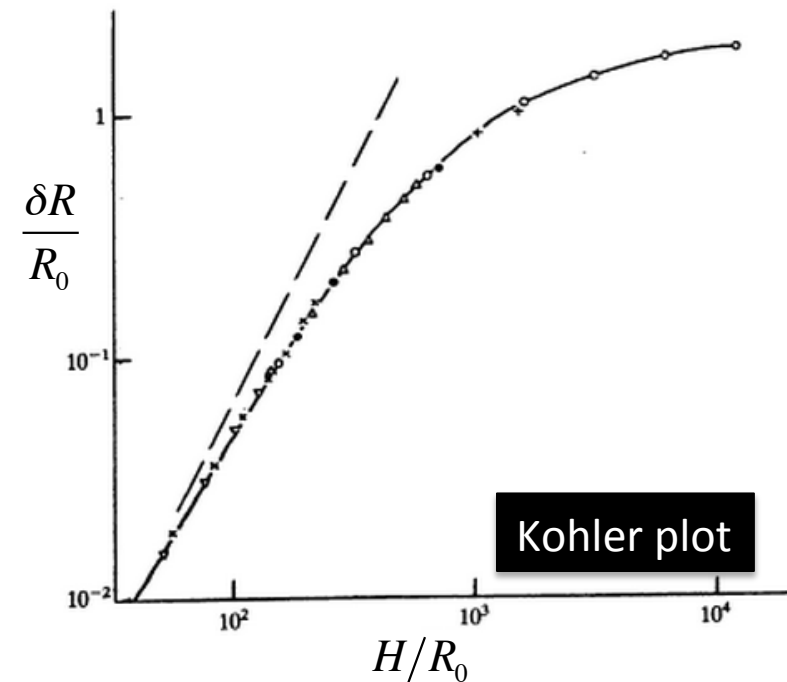
# Magnetoresistance and Kohler's rule

Magnetic field	$H \rightarrow cH$
Scattering rate	$\frac{1}{\tau} \rightarrow c \frac{1}{\tau}$
Resistivity	$\rho \rightarrow c\rho$

$$\longrightarrow \frac{\rho(H) - \rho_0}{\rho_0} = \frac{\delta\rho}{\rho_0} = F\left(\frac{H}{\rho_0}\right)$$

Indium with different impurity levels  
at the same temperature

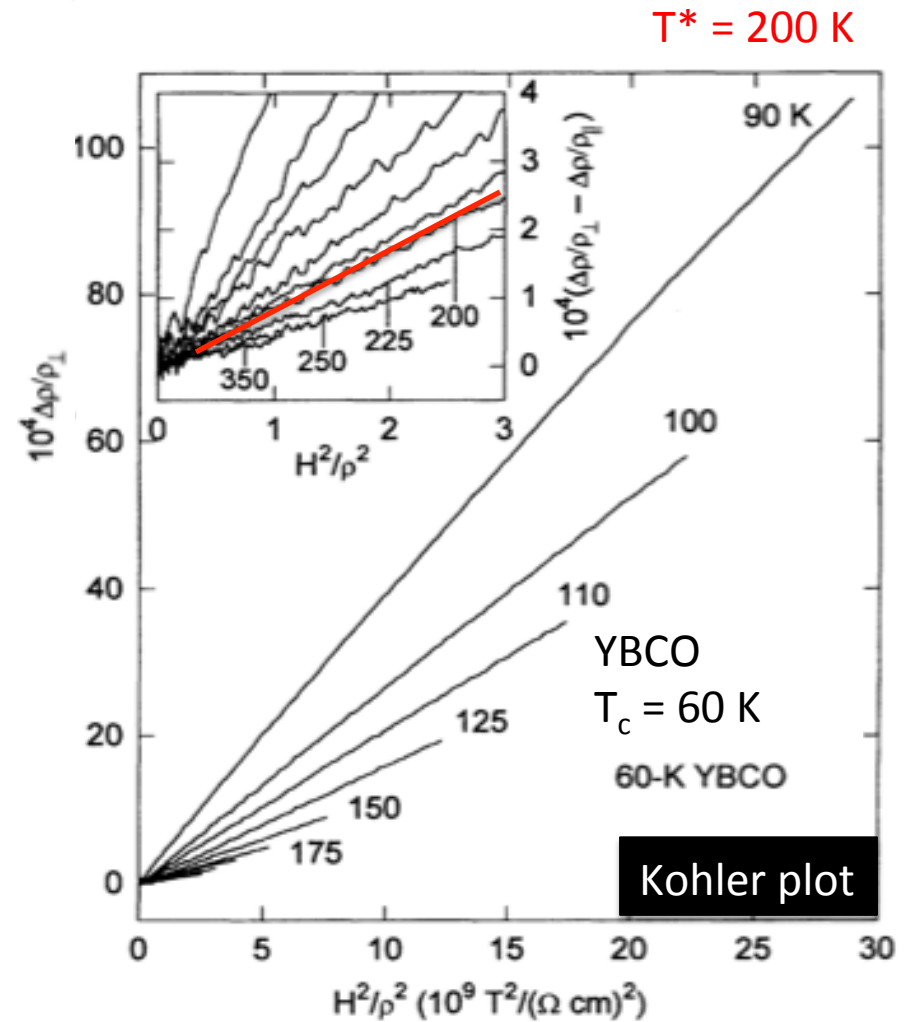
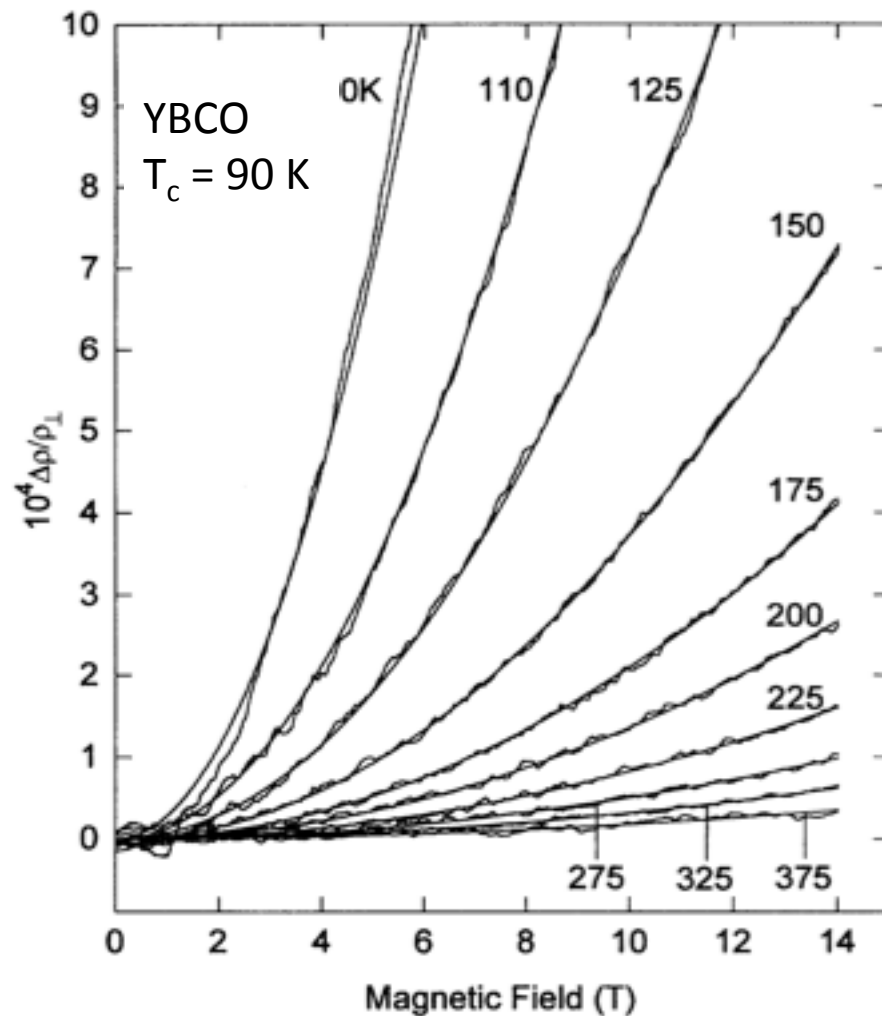
- Contingent on:
  - Only the amount of scattering changes
  - Single scattering mechanism
  - Constant number of carriers



J. L. Olsen *Electron transport in metals* (1962)

# Violation of Kohler's rule

- Violation of Boltzmann kinetics



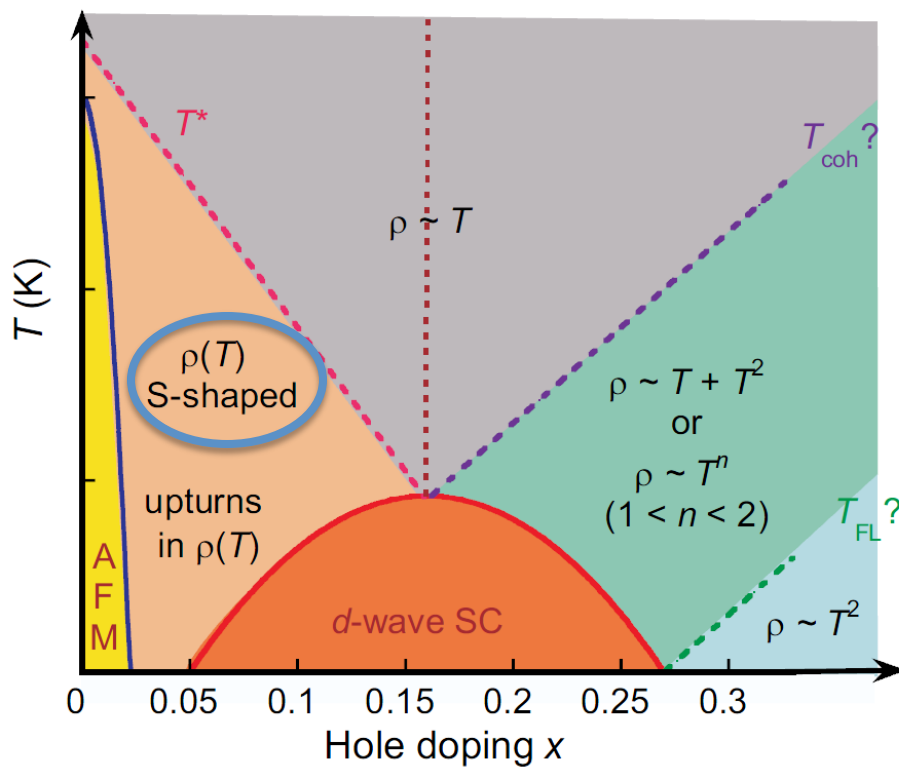
M. Harris *et al.* PRL (1995)

# Explanations for unconventional transport

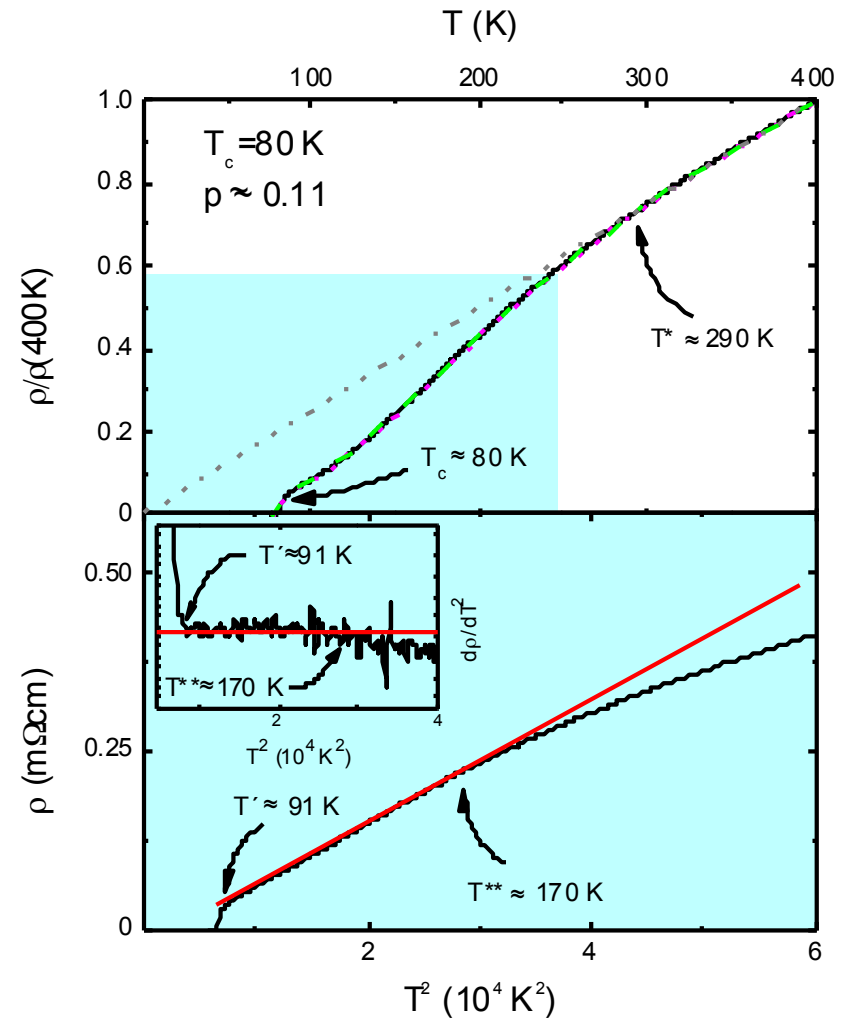
- Charge-spin separation (P. W. Anderson, Phys. Rev. Lett. **67**, 2092 (1991))
- Anisotropic Umklapp scattering (N. E. Hussey *et al.* PRL (1996))
- Coupling to a bosonic mode
  - Spin fluctuations (P. Monthoux and D. Pines, PRB (1994))
  - Charge fluctuations (C. Castellani *et al.* PRL (1995))

# Electrical transport in the pseudogap

- Electrical transport is ill-defined in the pseudogap phase
- Fluctuations of superconductivity or nearby ordered states



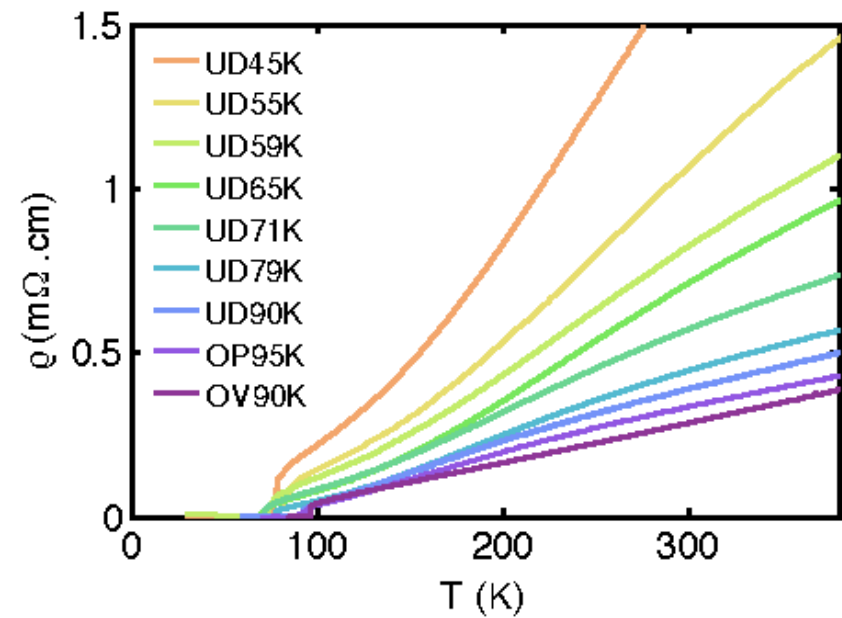
N. J. Hussey Phys.: Condens. Matter. (2008)



N. Barišić, MKC *et al.* PNAS (2013)

# Electrical transport in the pseudogap

- Clear separation of a linear and quadratic regimes

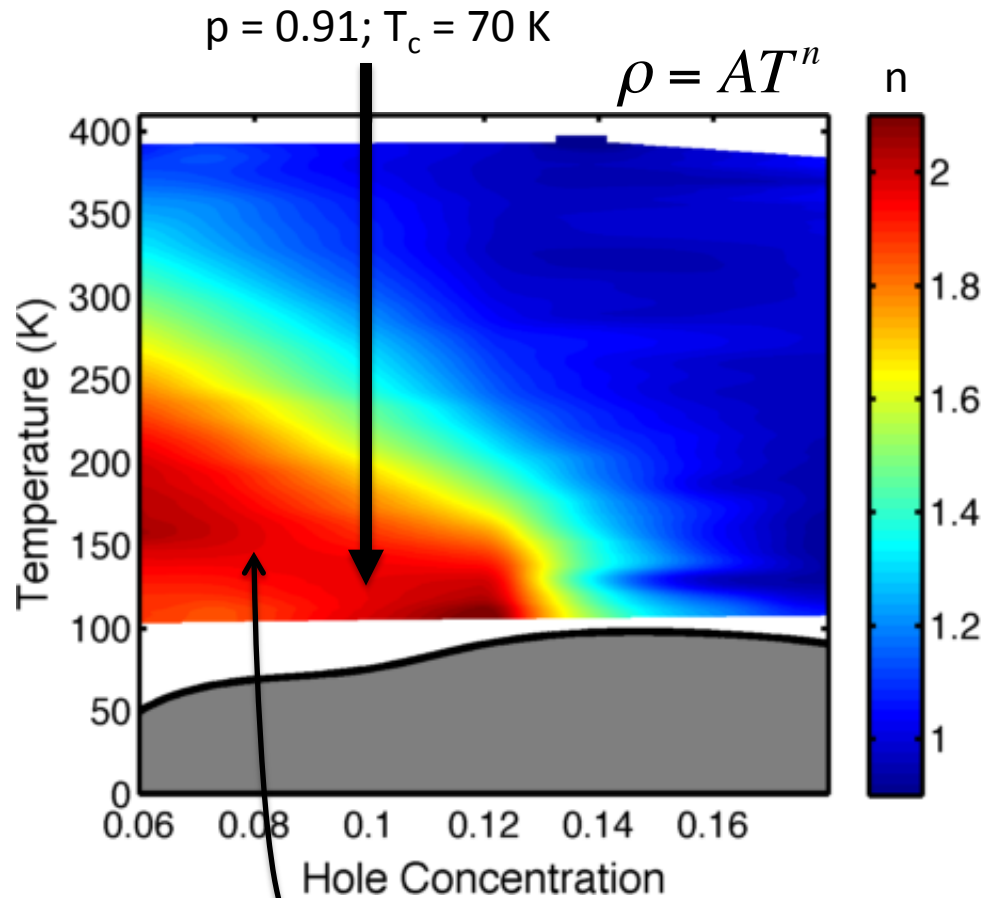


Boltzmann kinetics?

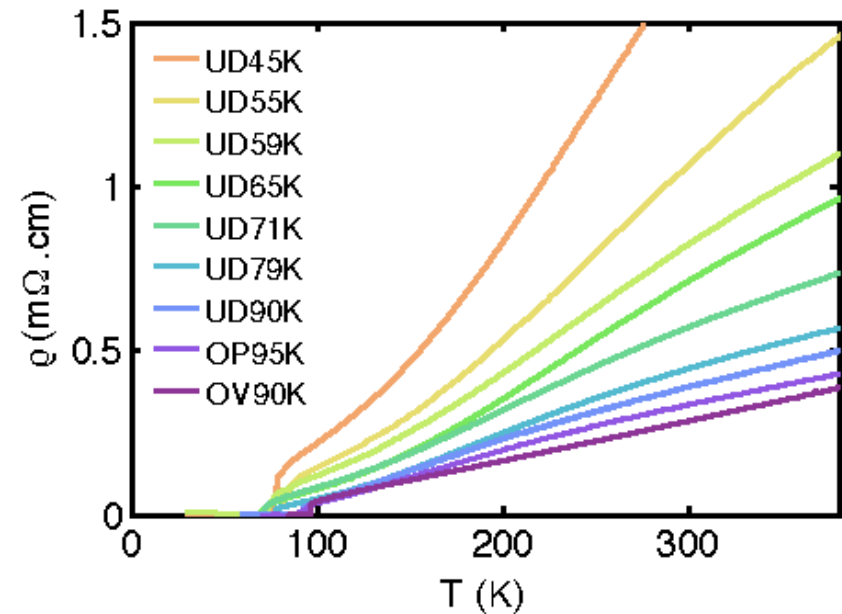
M .K. Chan *et al.* unpublished

# Magnetoresistance

- Clear separation of a linear and quadratic regimes



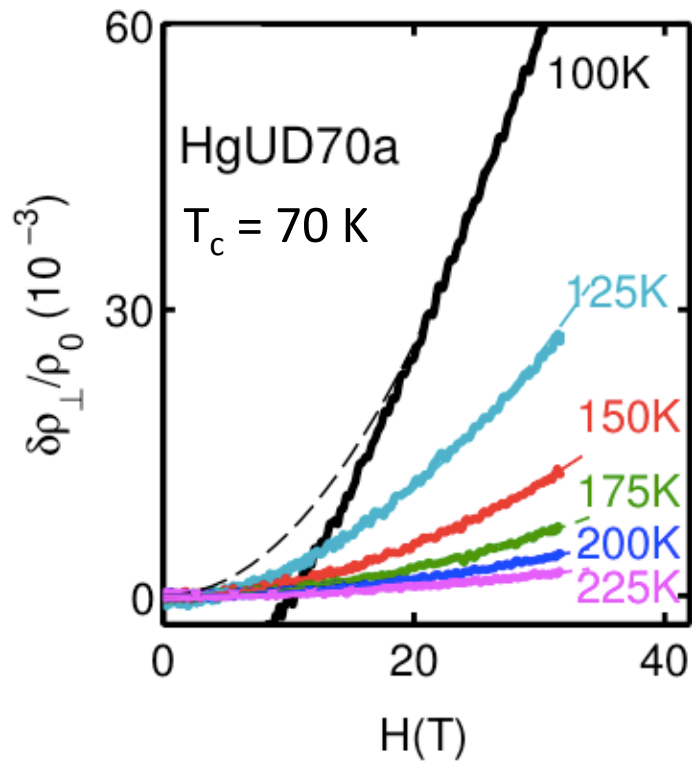
Boltzmann kinetics?



M .K. Chan *et al.* unpublished

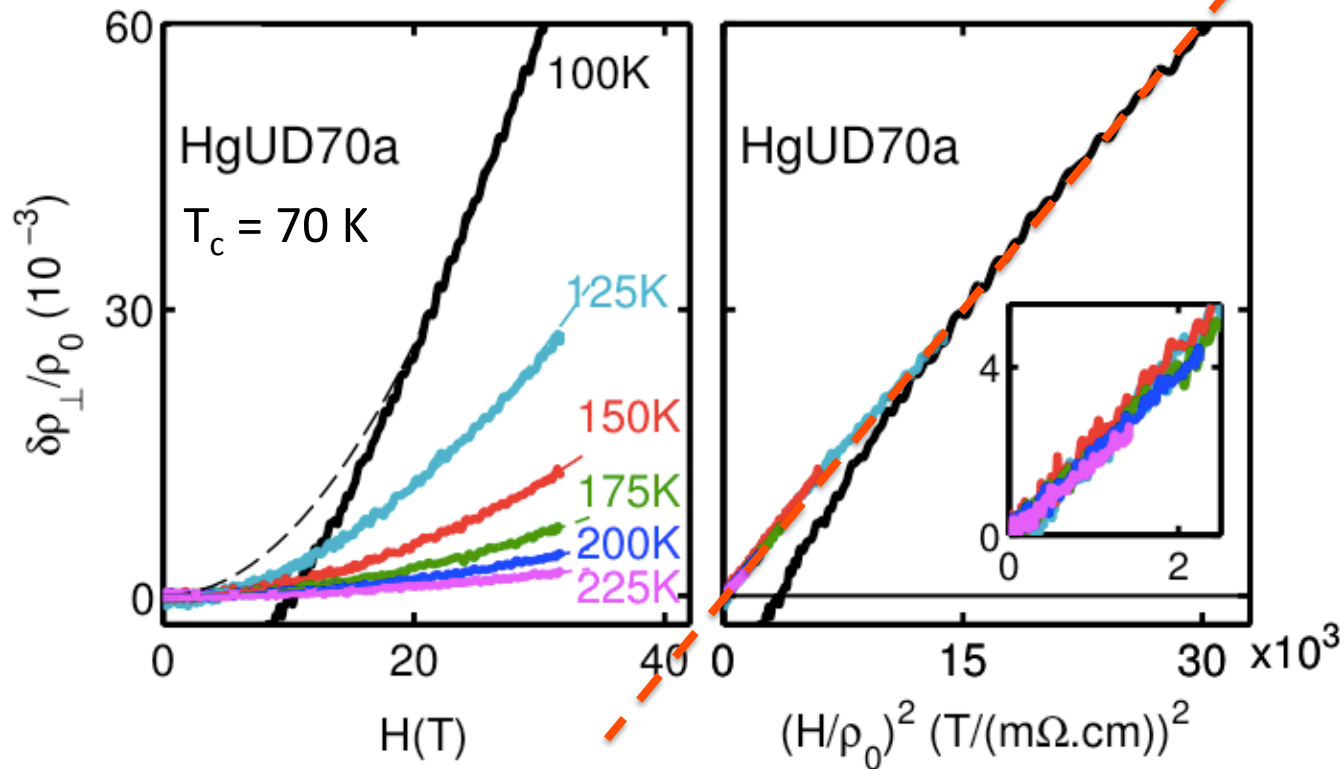
# Validity of Kohler's rule

- Pulsed field measurements at LCNMI-Toulouse. 30 T.  $j \parallel ab$ ;  $H \parallel c$
- $\delta\rho/\rho_0 = a' H^2$



# Validity of Kohler's rule

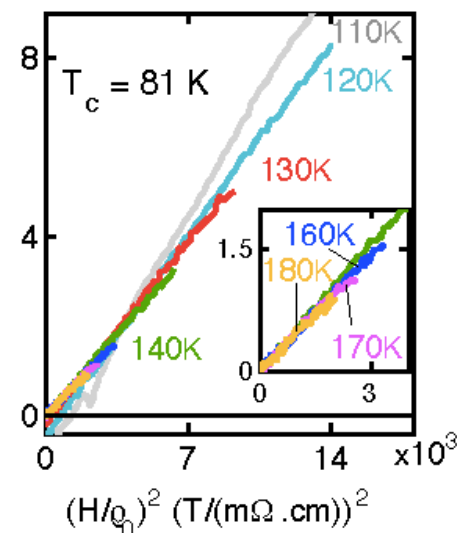
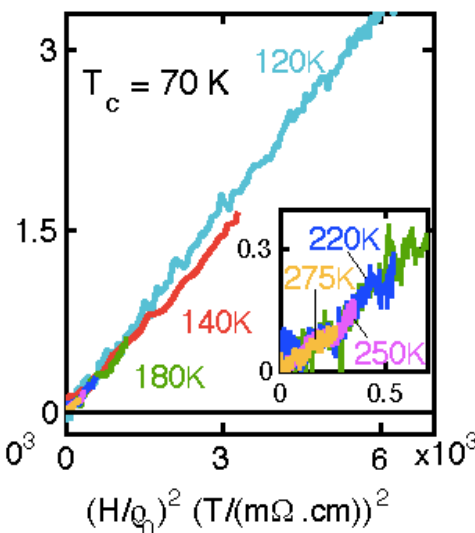
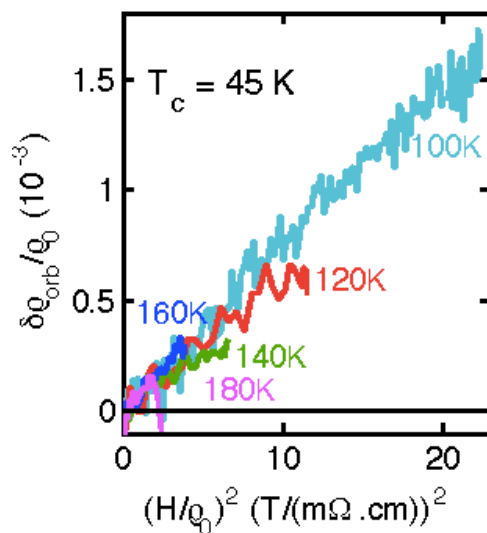
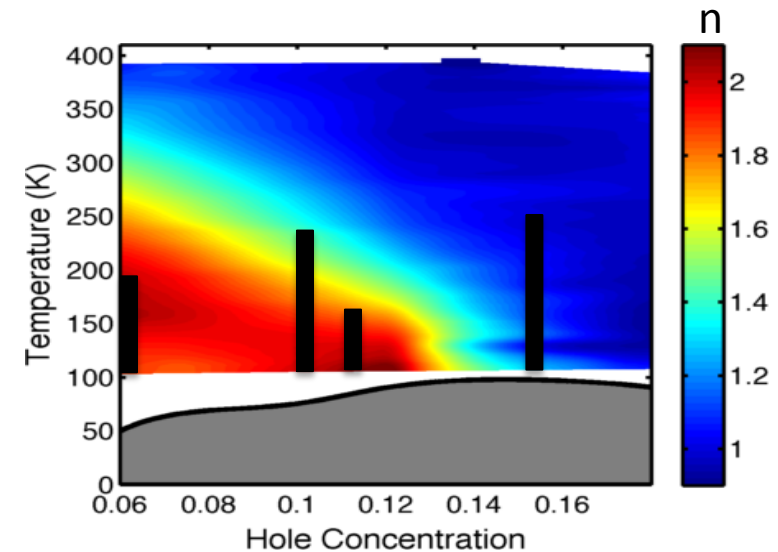
- Pulsed field measurements at LCNMI-Toulouse. 30 T.  $j \parallel ab$ ;  $H \parallel c$
- $\delta\rho/\rho_0 = a' H^2$
- **Kohler's rule is valid in the pseudogap**



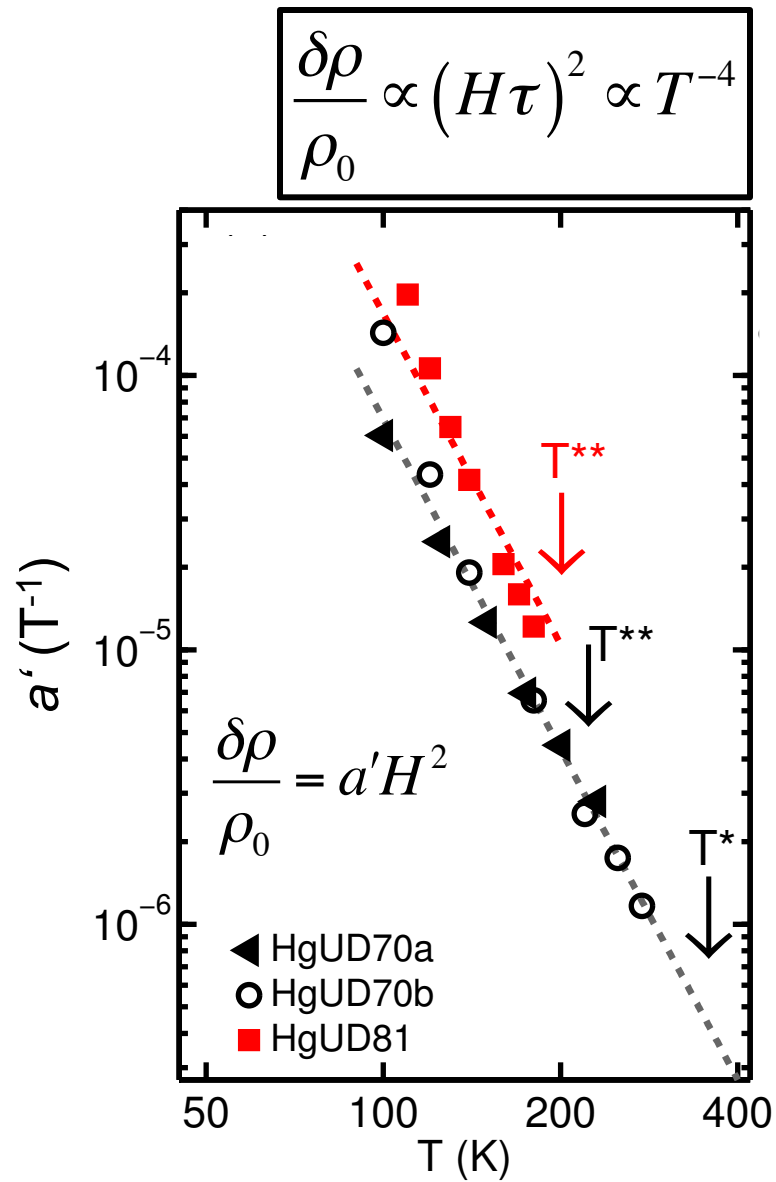


# Doping dependence of Kohler's rule

- In-house measurements :  $H = 9$  Tesla
- Kohler's rule valid throughout pseudogap regime where  $\rho \sim T^2$
- Kohler's rule broken in the strange-metal regime:  $\rho \sim T^1$



# Temperature dependence of MR

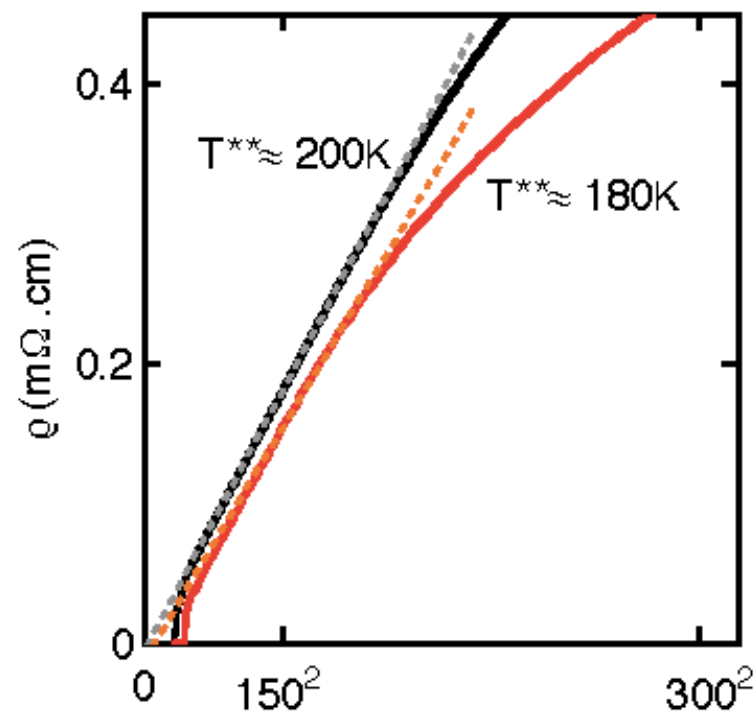
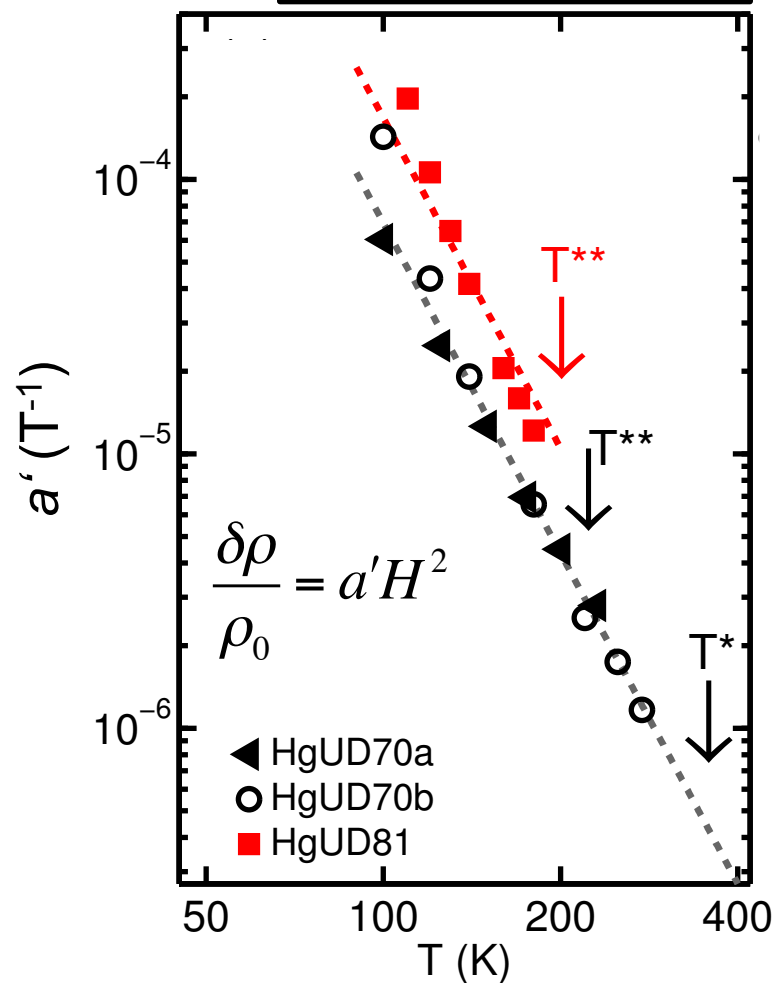


# Temperature dependence of MR

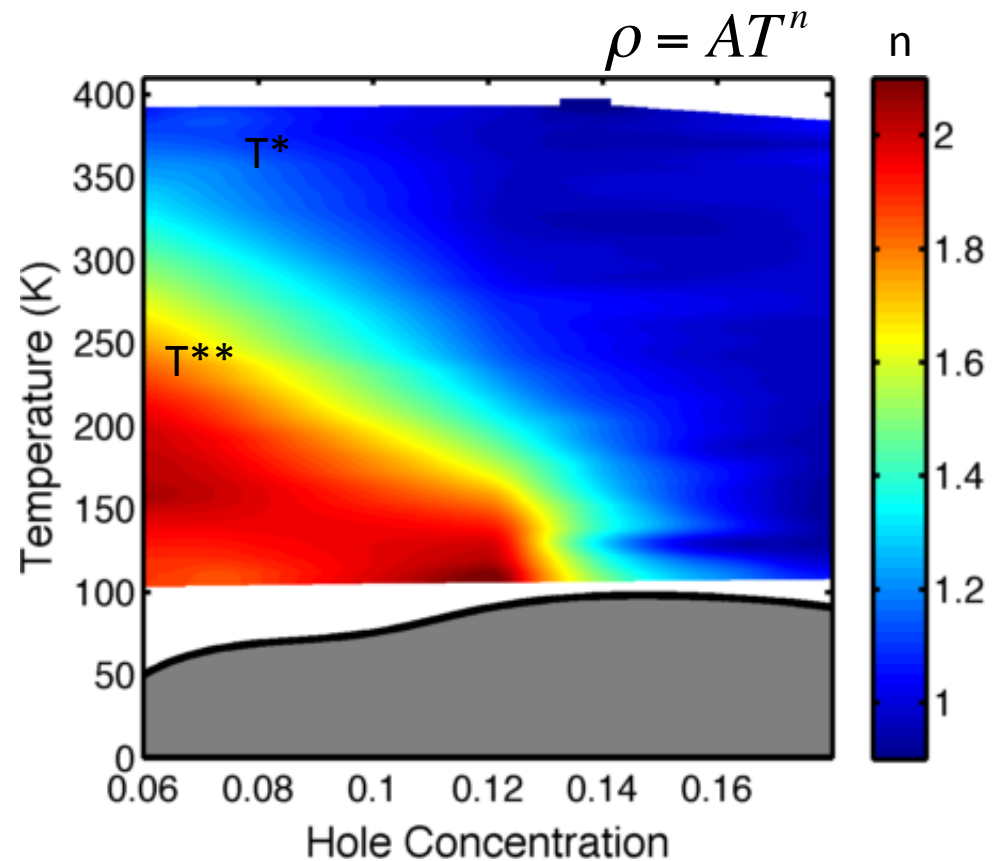
$$\frac{\delta\rho}{\rho_0} \propto (H\tau)^2 \propto T^{-4}$$

$$\rho \propto (1/\tau) \propto T^2$$

Kohler's rule

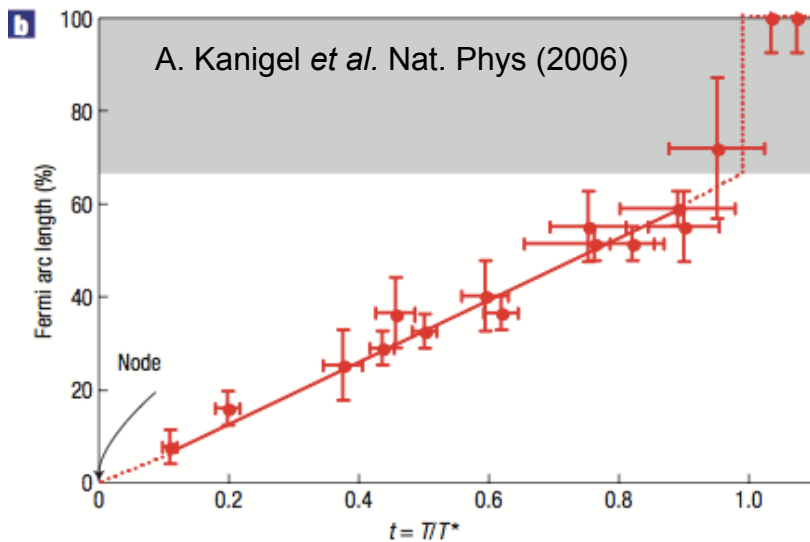
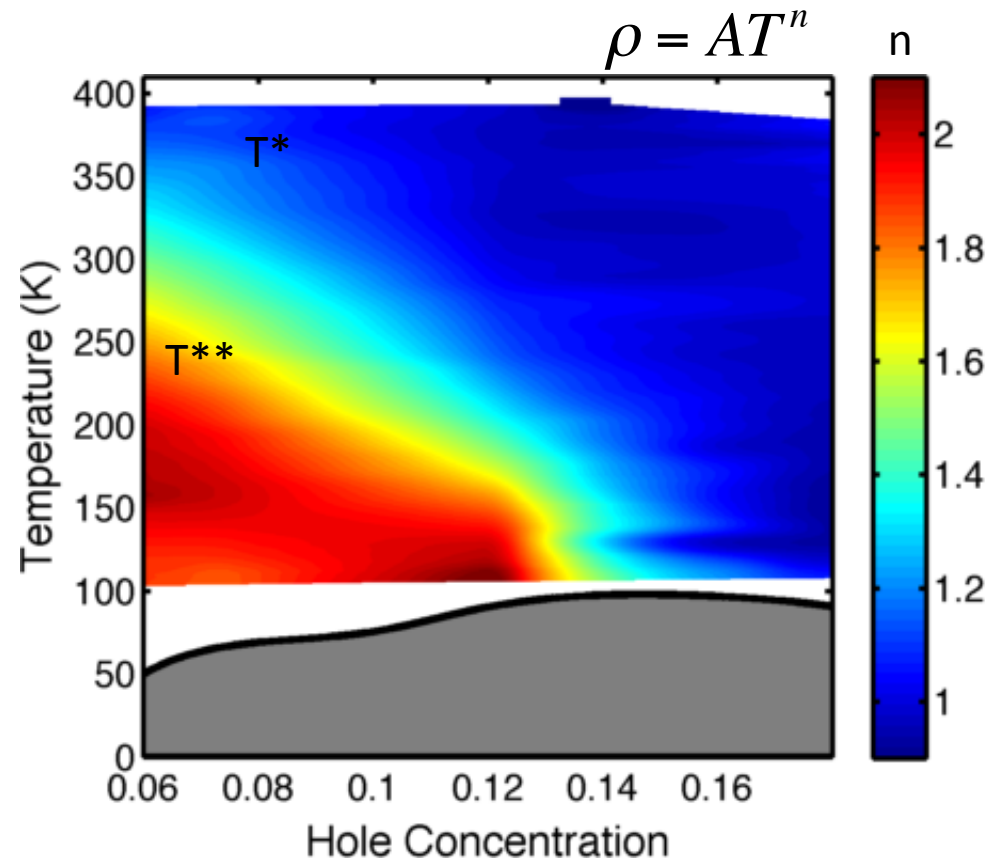
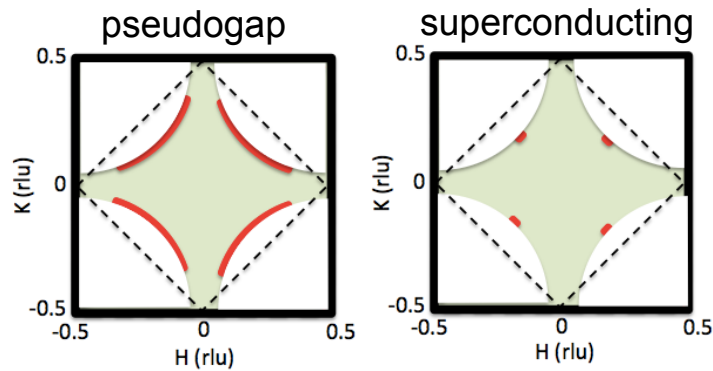


# Implications



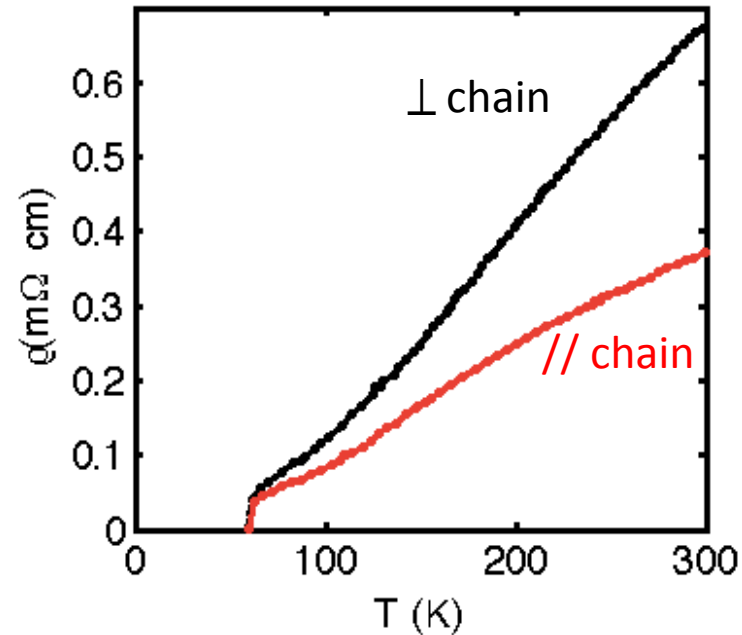
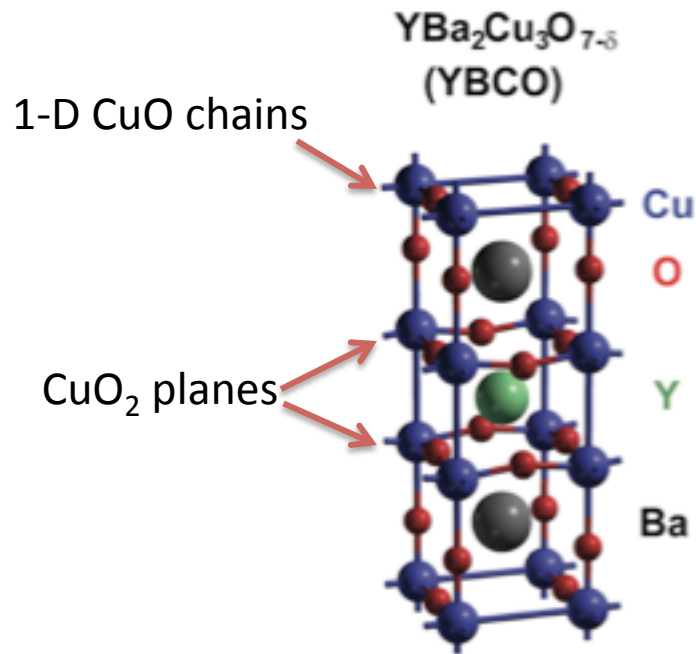
- Only **one** Fermi-liquid-like scattering rate
- Constant carrier density
- Stable Fermi surface between  $T_c$  and  $T^{**}$

# Implications



- Pseudogap is not fluctuating superconductivity
- Stable Fermi arc in the pseudogap supported by more recent ARPES results in optimally doped Bi-compounds: T. Kondo *et al.* PRL (2013)

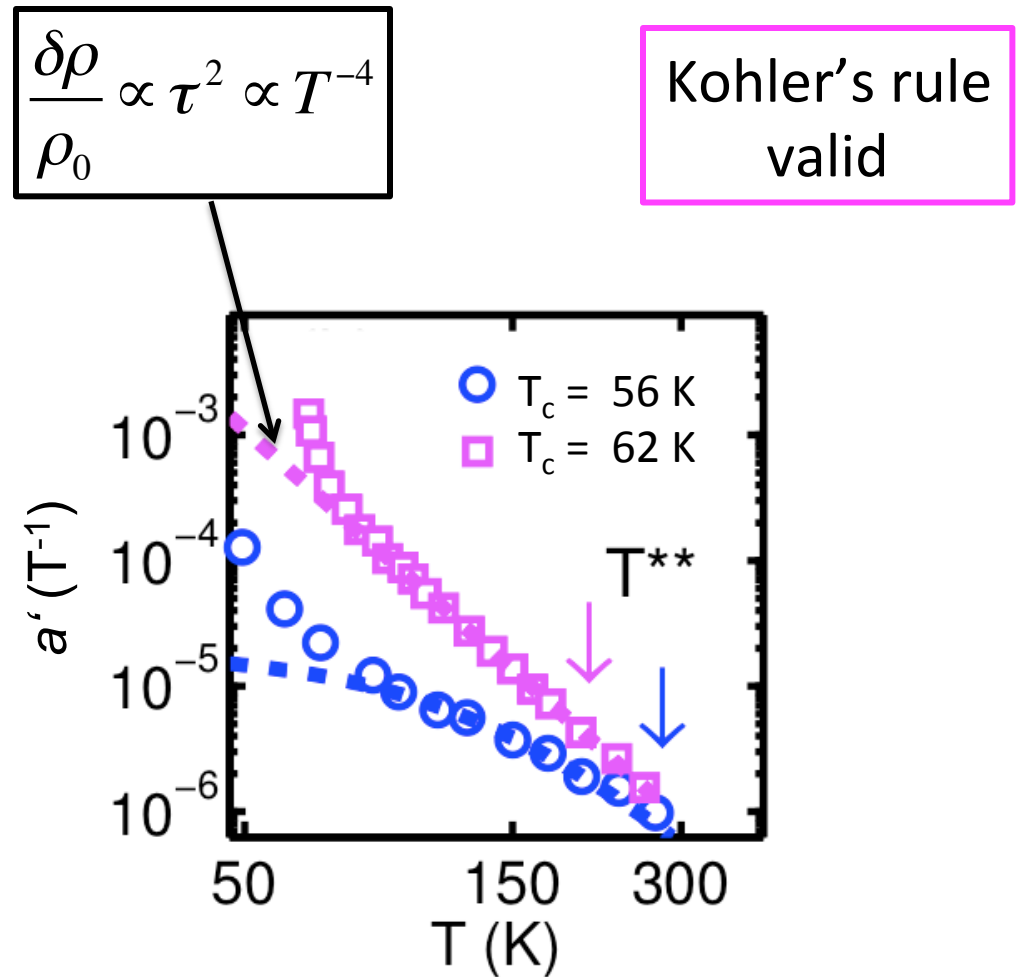
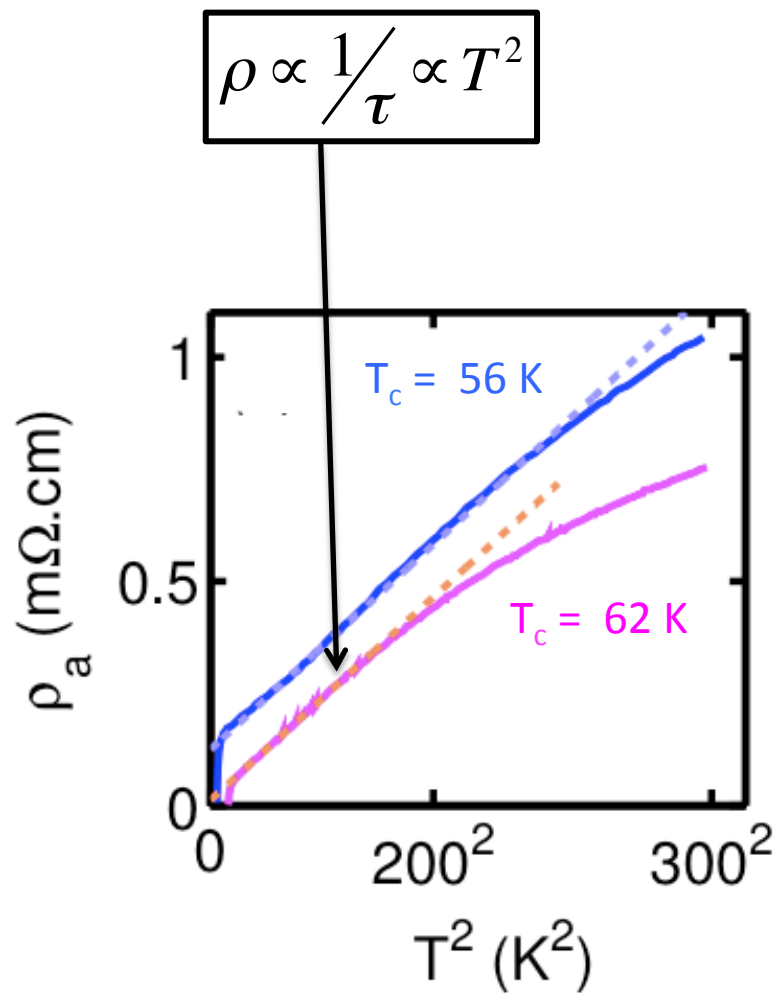
# The case of YBCO



Adapted from Y. Ando *et al.* PRL (2002)

- CuO chains contribute to the electrical transport : measurements in twinned crystals naturally violate Kohler's rule
- Chains shouldn't contribute to the electrical transport transverse to them

## Transport perpendicular to CuO chains



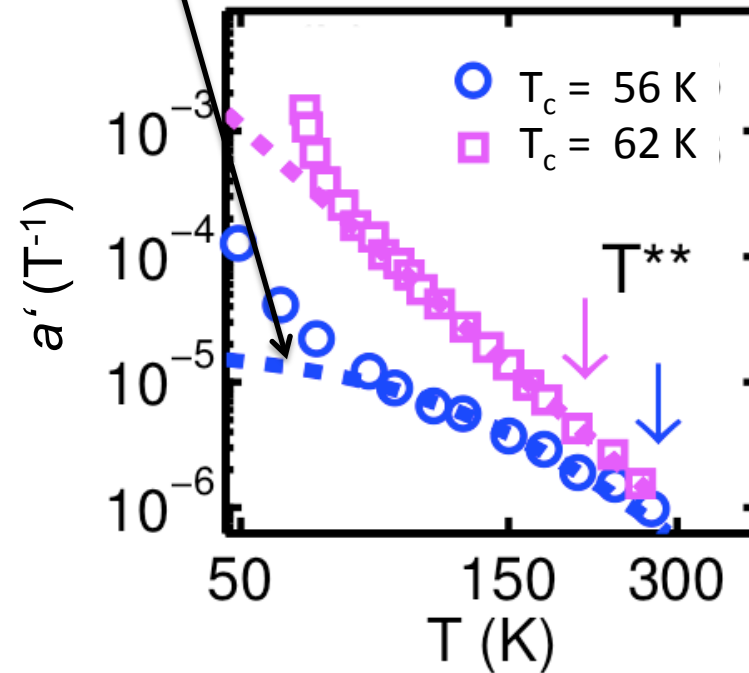
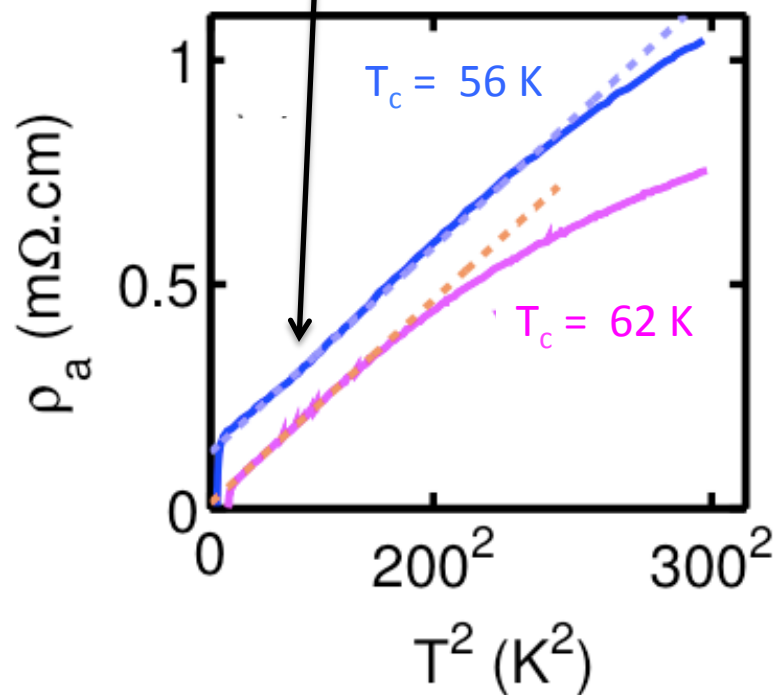
Adapted from Y. Ando *et al.* PRL (2002)

# Transport perpendicular to CuO chains

$$\rho \propto \frac{1}{\tau} = b + cT^2$$

$$\frac{\delta\rho}{\rho_0} \propto \tau^2 = \frac{1}{(b + cT^2)^2}$$

Kohler's rule  
not valid

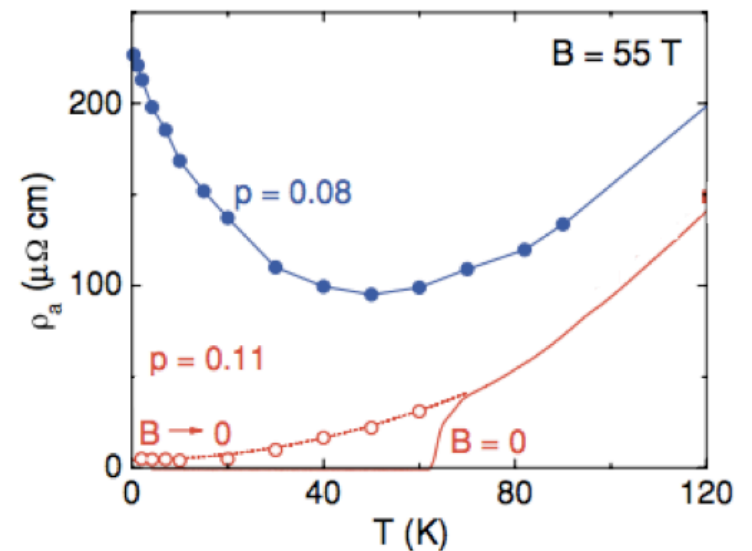
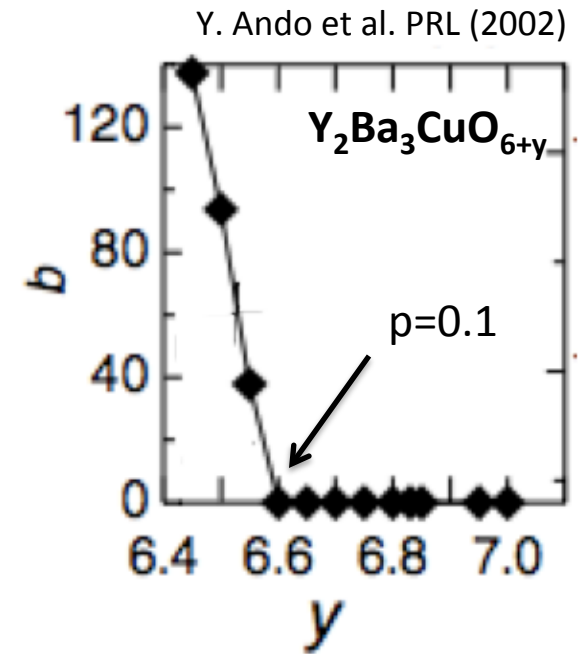
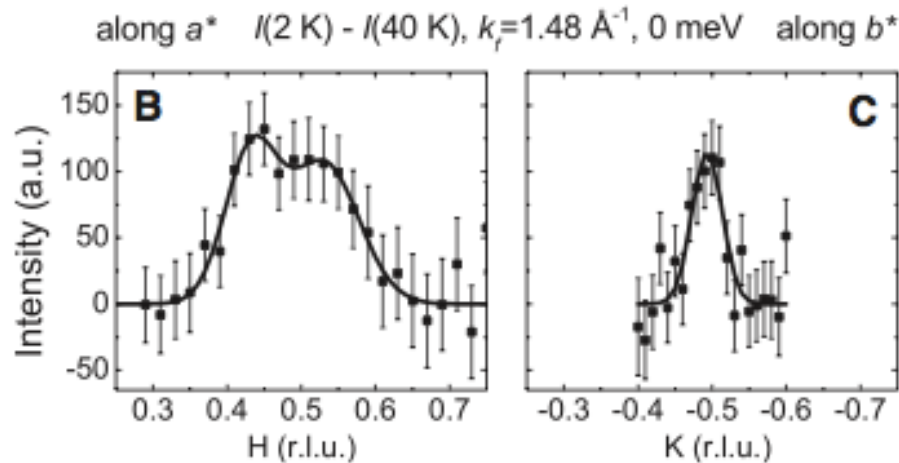
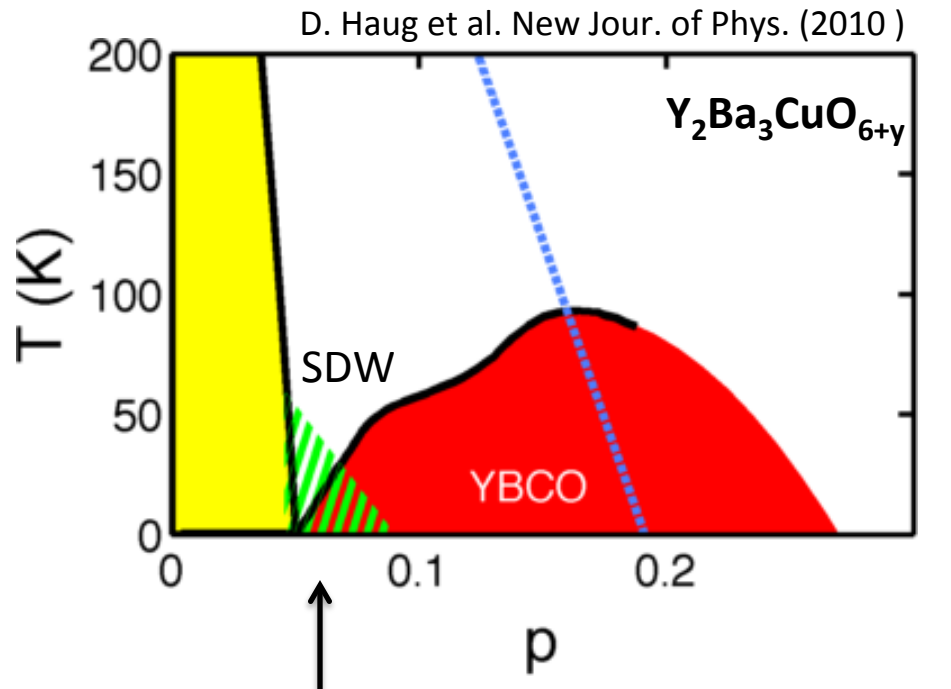


Adapted from Y. Ando *et al.* PRL (2002)



# Magnetic phase diagram and residual resistivity

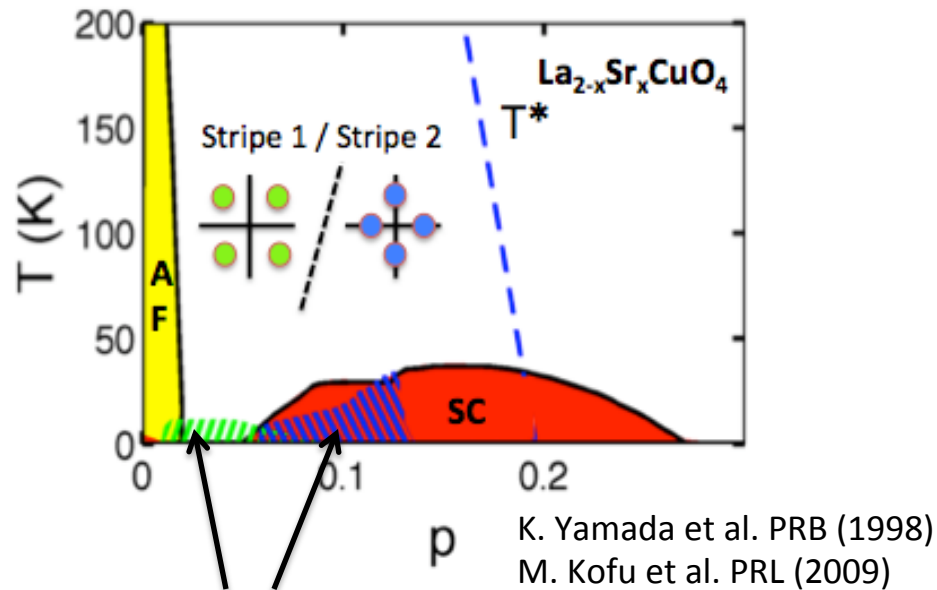
## ○ 3D AF → incommensurate SDW



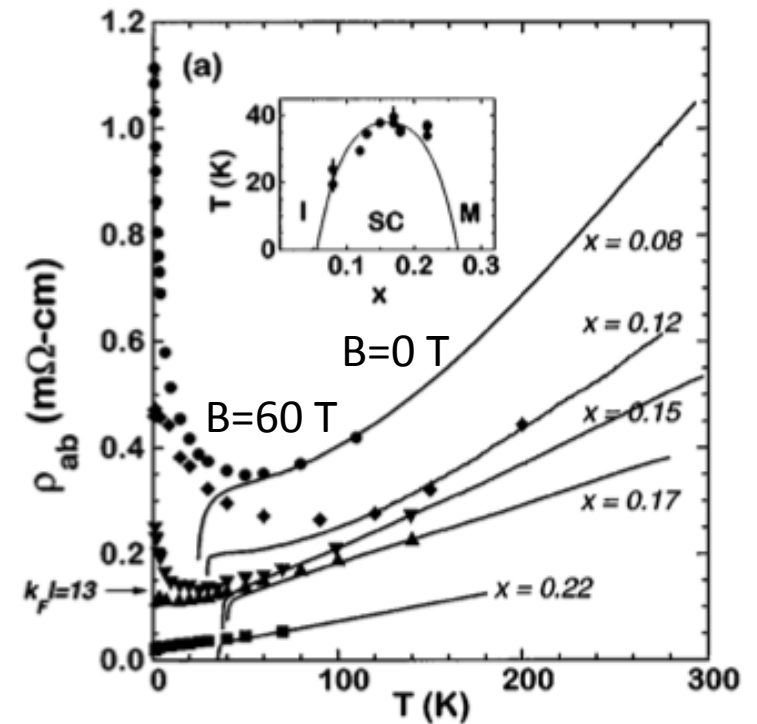
D. LeBoeuf et al. PRB (2011)

# Magnetic phase diagram and residual resistivity

- 3D AF → diagonal incommensurate → parallel incommensurate



Local moment physics

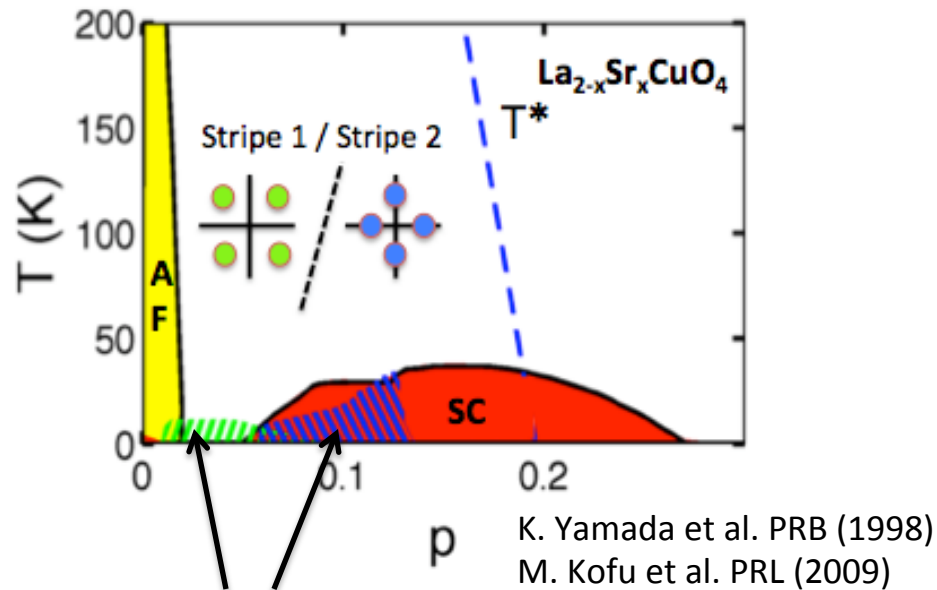


G. S. Boebinger et al. PRL (1996)

- Large residual resistivity is an indicator of carrier localization at low temperatures and quasielastic magnetic ordering

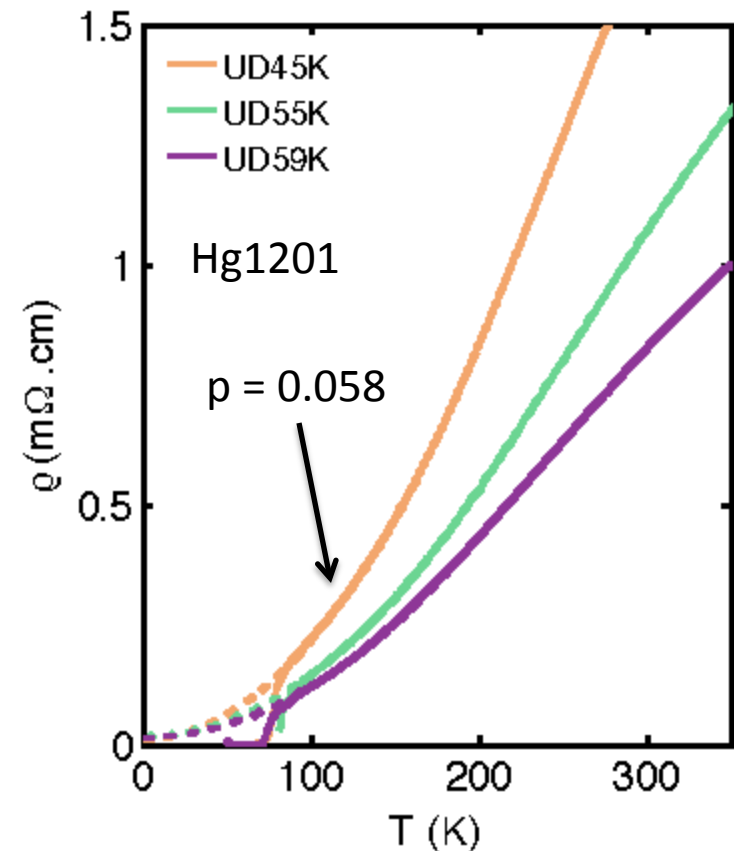
# Magnetic phase diagram and residual resistivity

- 3D AF → diagonal incommensurate → parallel incommensurate



Local moment physics

- $p_{\text{MI}}$  (LSCO)  $\approx .17$
- $p_{\text{MI}}$  (YBCO)  $\approx .10$
- $p_{\text{MI}}$  (Hg1201)  $< 0.058$
- Quantum oscillations in  $p=0.09$  (N. Barišić *et al.* Nature Phys. (2013))



## Part 2: Inelastic Neutron scattering

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## Part 2: Inelastic Neutron scattering



Yang Tang  
(UMN)

### LLB, CEA-Saclay, France

Philippe Bourges  
Yvan Sidis  
Lucille Mangin-Thro

### ILL, Grenoble, France

Paul Steffens

### FRM-II, Garching, Germany

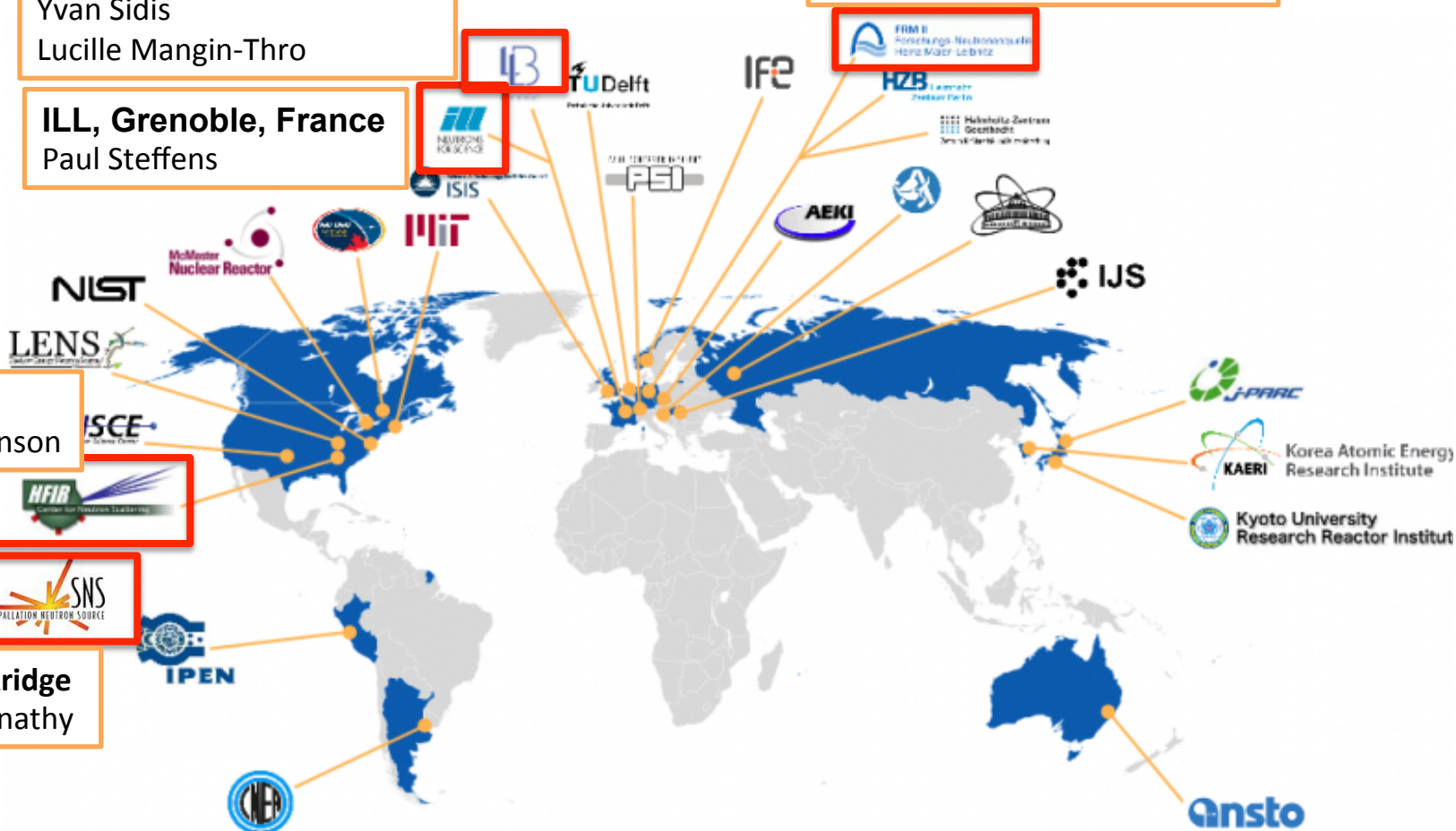
Jitae Park

### ORNL, Oakridge

Andrew Christianson

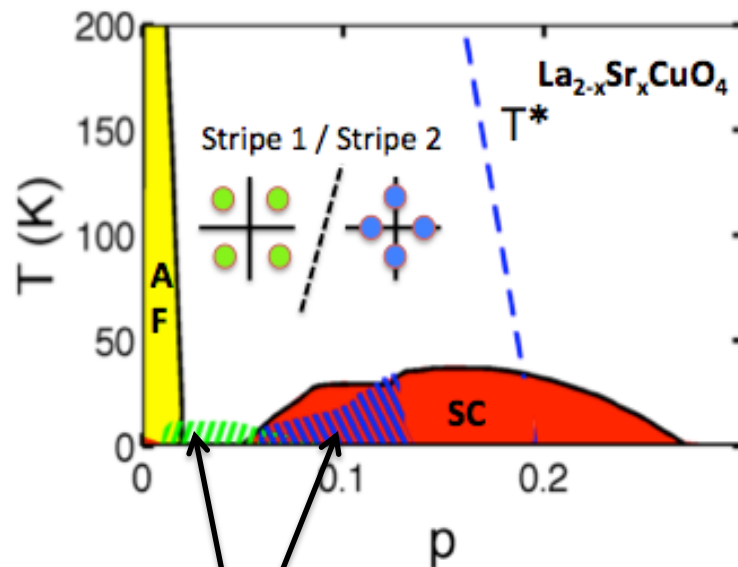
### ORNL, Oakridge

Doug Abernathy

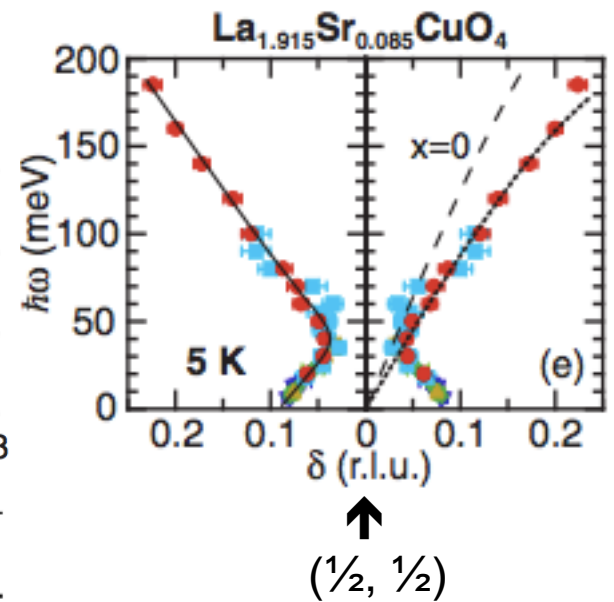
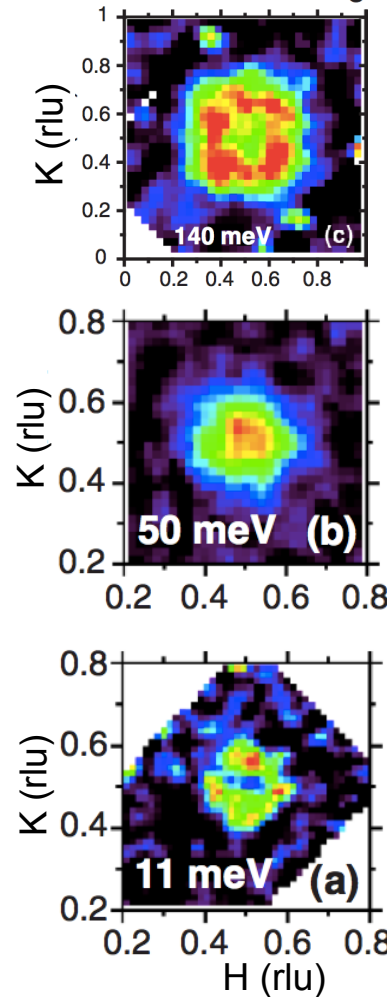


# Magnetic phase diagram

- 3D AF  $\rightarrow$  diagonal incommensurate  $\rightarrow$  parallel incommensurate

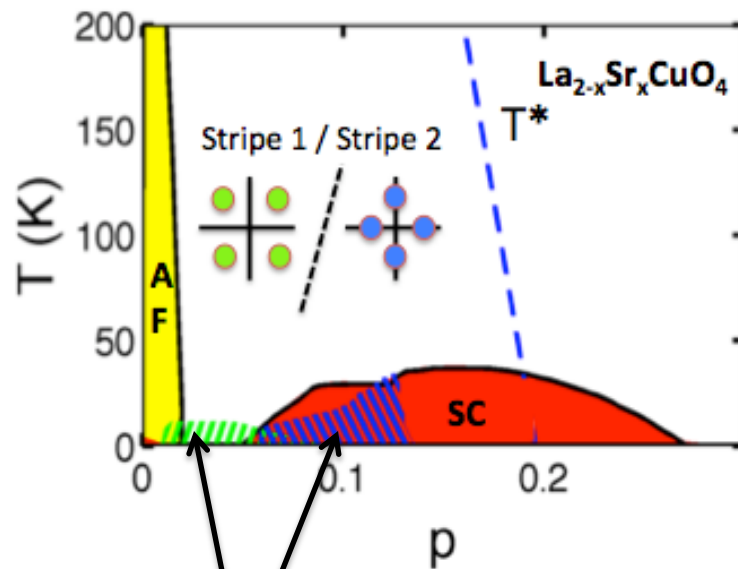


Local moment physics

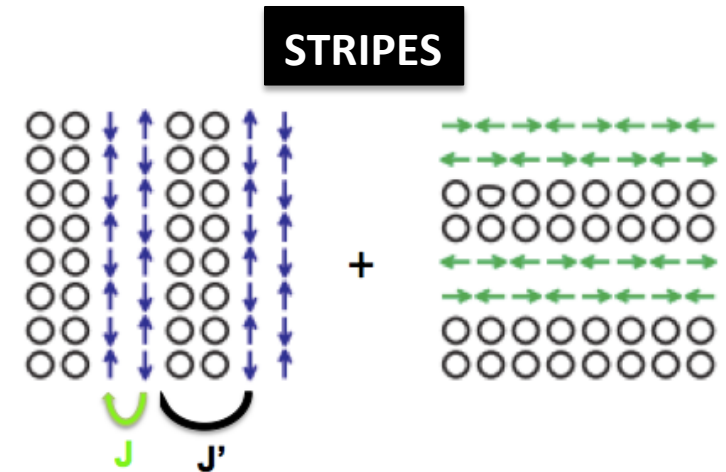


# Magnetic phase diagram

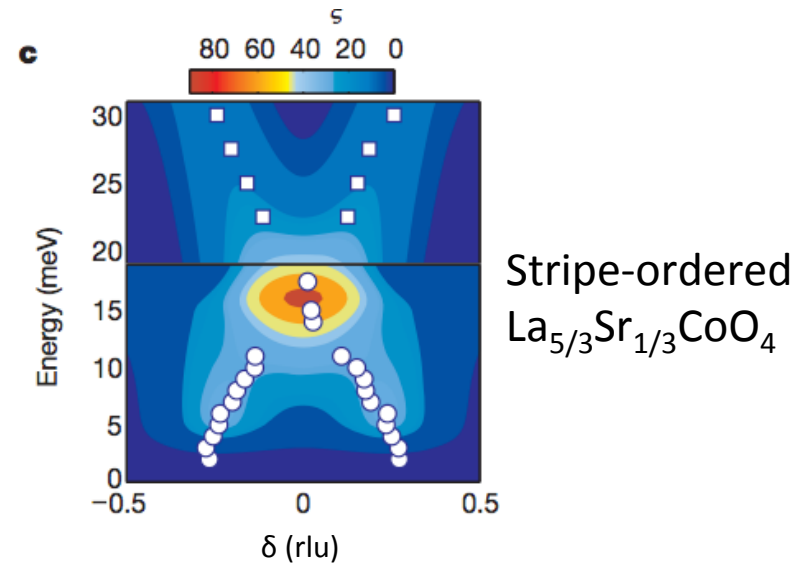
- 3D AF → diagonal incommensurate → parallel incommensurate



Local moment physics



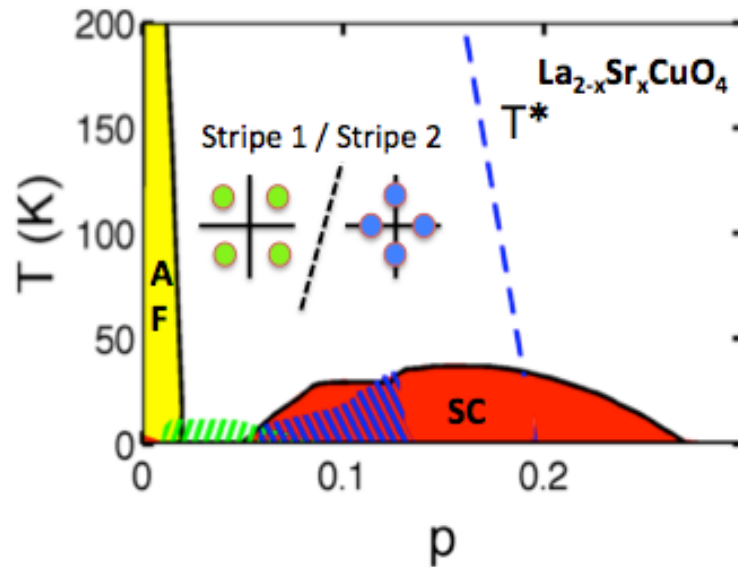
J. M. Tranquada et al. Nature 429 534 (2004)



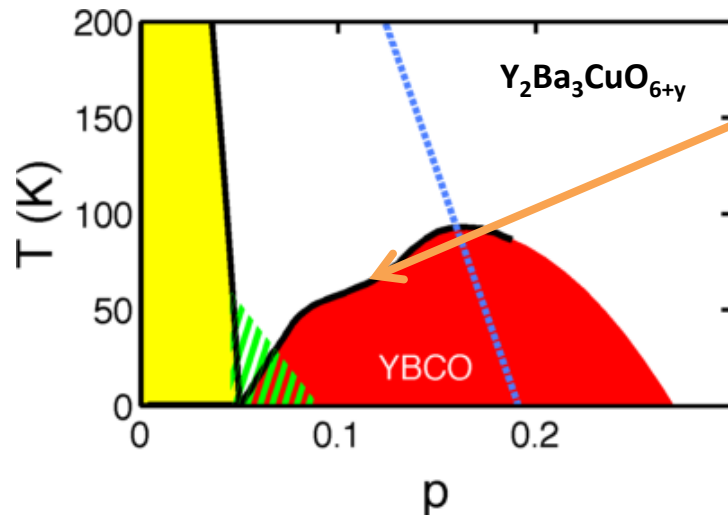
A. T. Boothroyd et al. Nature (2011)

# Magnetic phase diagram

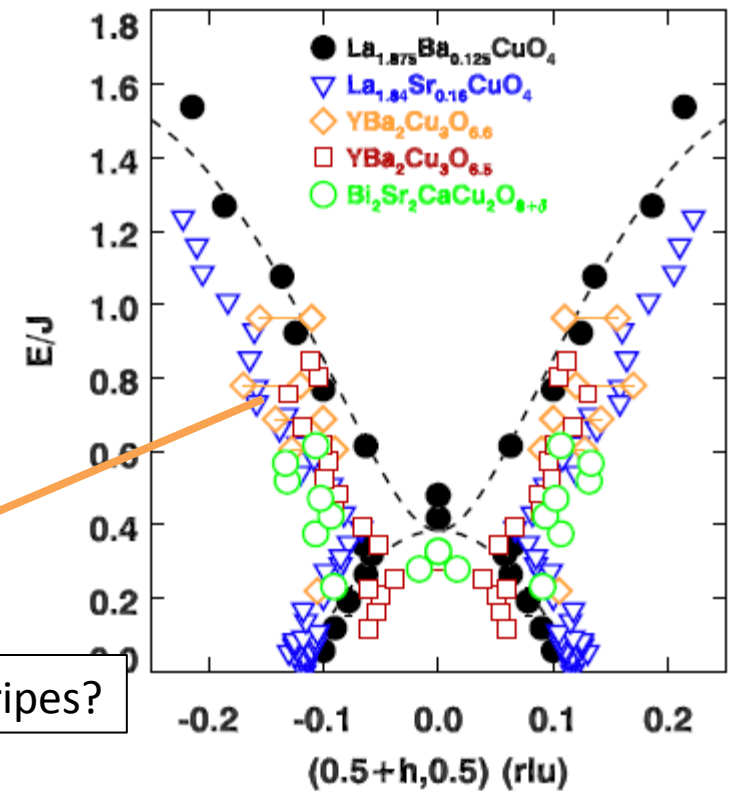
- 3D AF → diagonal incommensurate → parallel incommensurate



- 3D AF → anisotropic incommensurate → ?



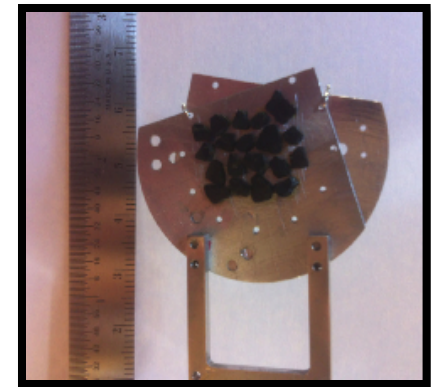
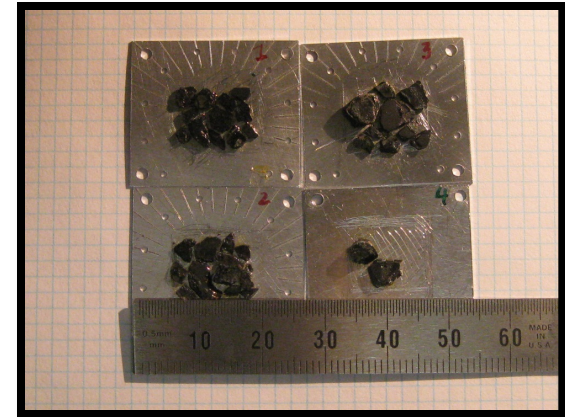
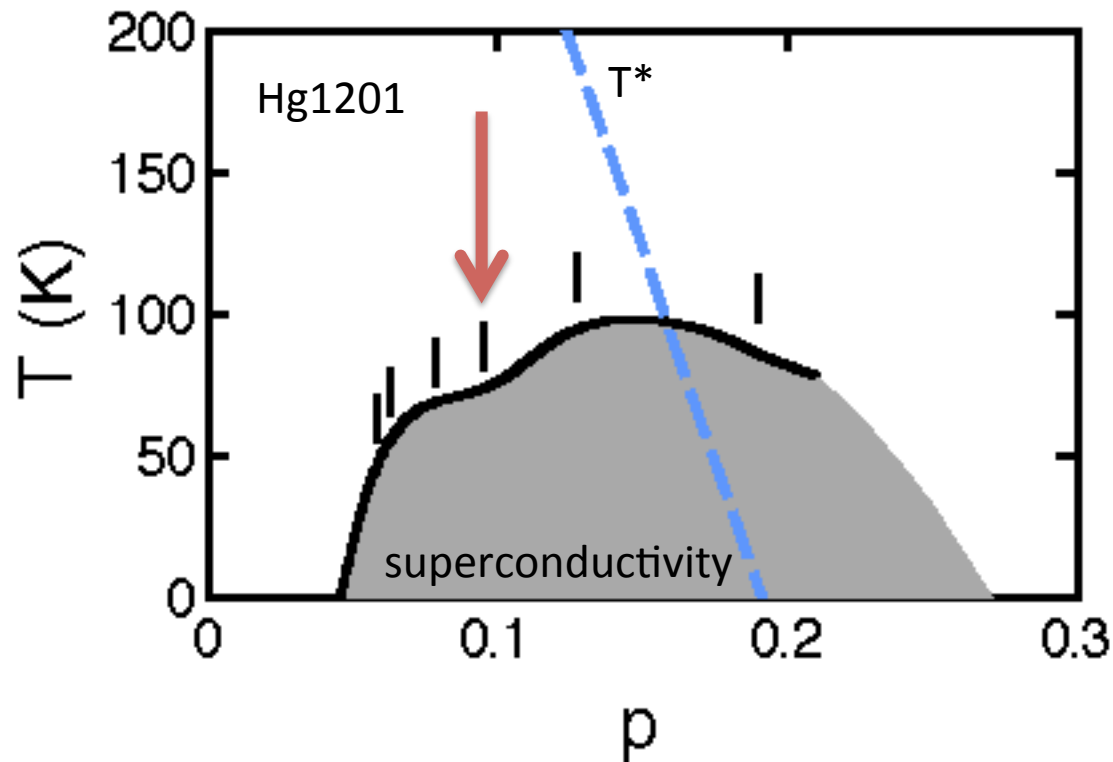
Fluctuating stripes?



M. Fujita *et al.* Physica C (2012)



# Preparing a neutron sample



Each neutron sample:

- ~ 40 growths
- ~30 crystals, ranging from 20 mg – 100 mg each
- 2 – 3 months anneal
- 2 g total mass

# Antiferromagnetic Fluctuations

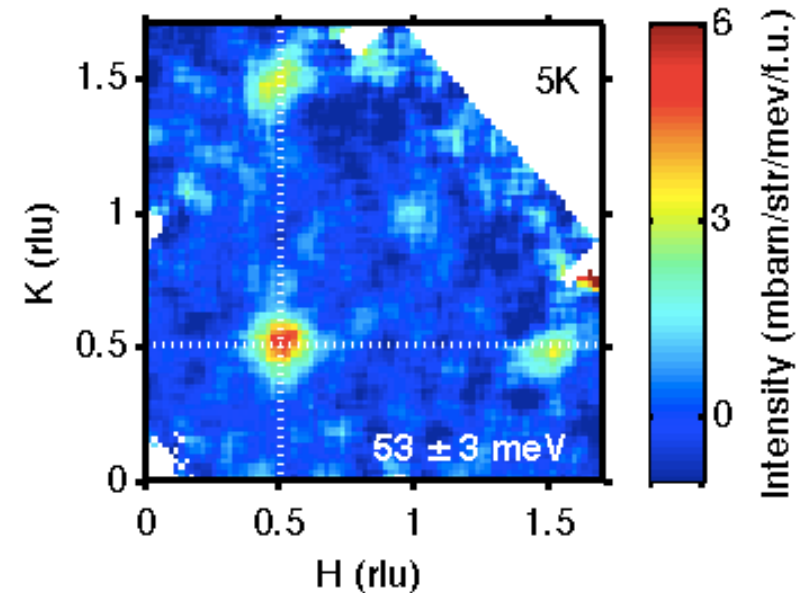
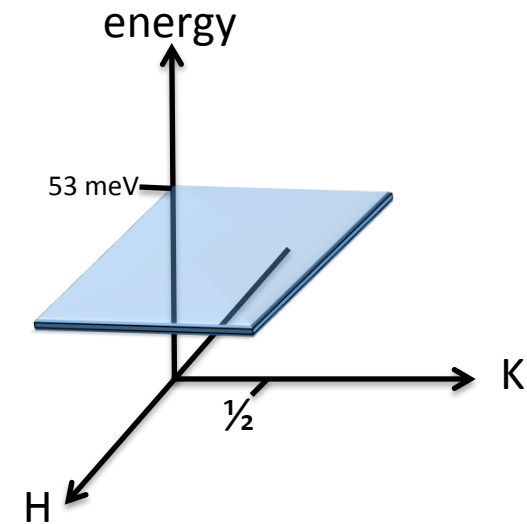
$$\frac{d^2\sigma}{d\Omega dE} = \frac{2(\gamma r_e)^2 k_f}{\pi g^2 \mu_B^2 k_i} |F(\mathbf{Q})|^2 \frac{\chi''(\mathbf{Q}, \omega)}{1 - \exp(-\omega/k_B T)}$$

$|F(\mathbf{Q})|$  = Form Factor

$\frac{1}{1 - \exp(-\omega/k_B T)}$  = Bose factor

$\chi''(\mathbf{Q}, \omega)$  = Imaginary magnetic susceptibility

$k_i, k_f$  = incident and final momenta



# Antiferromagnetic Fluctuations

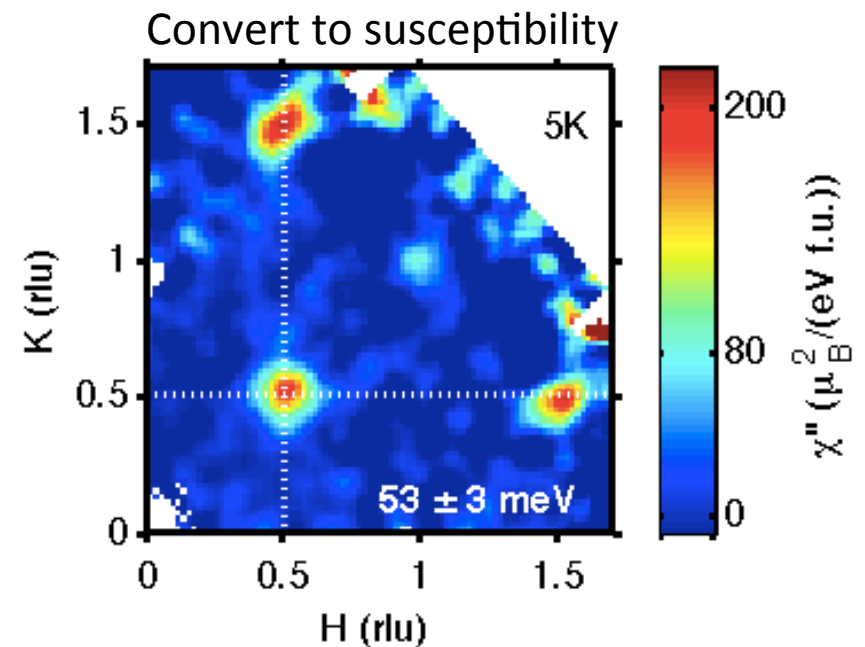
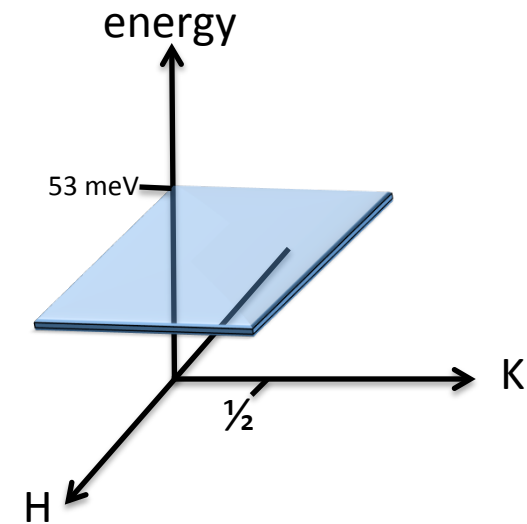
$$\frac{d^2\sigma}{d\Omega dE} = \frac{2(\gamma r_e)^2 k_f}{\pi g^2 \mu_B^2 k_i} |F(\mathbf{Q})|^2 \frac{\chi''(\mathbf{Q}, \omega)}{1 - \exp(-\omega/k_B T)}$$

$|F(\mathbf{Q})|$  = Form Factor

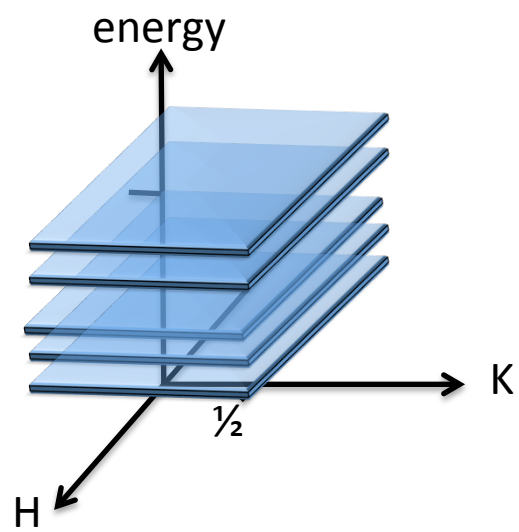
$\frac{1}{1 - \exp(-\omega/k_B T)}$  = Bose factor

$\chi''(\mathbf{Q}, \omega)$  = Imaginary magnetic susceptibility

$k_i, k_f$  = incident and final momenta

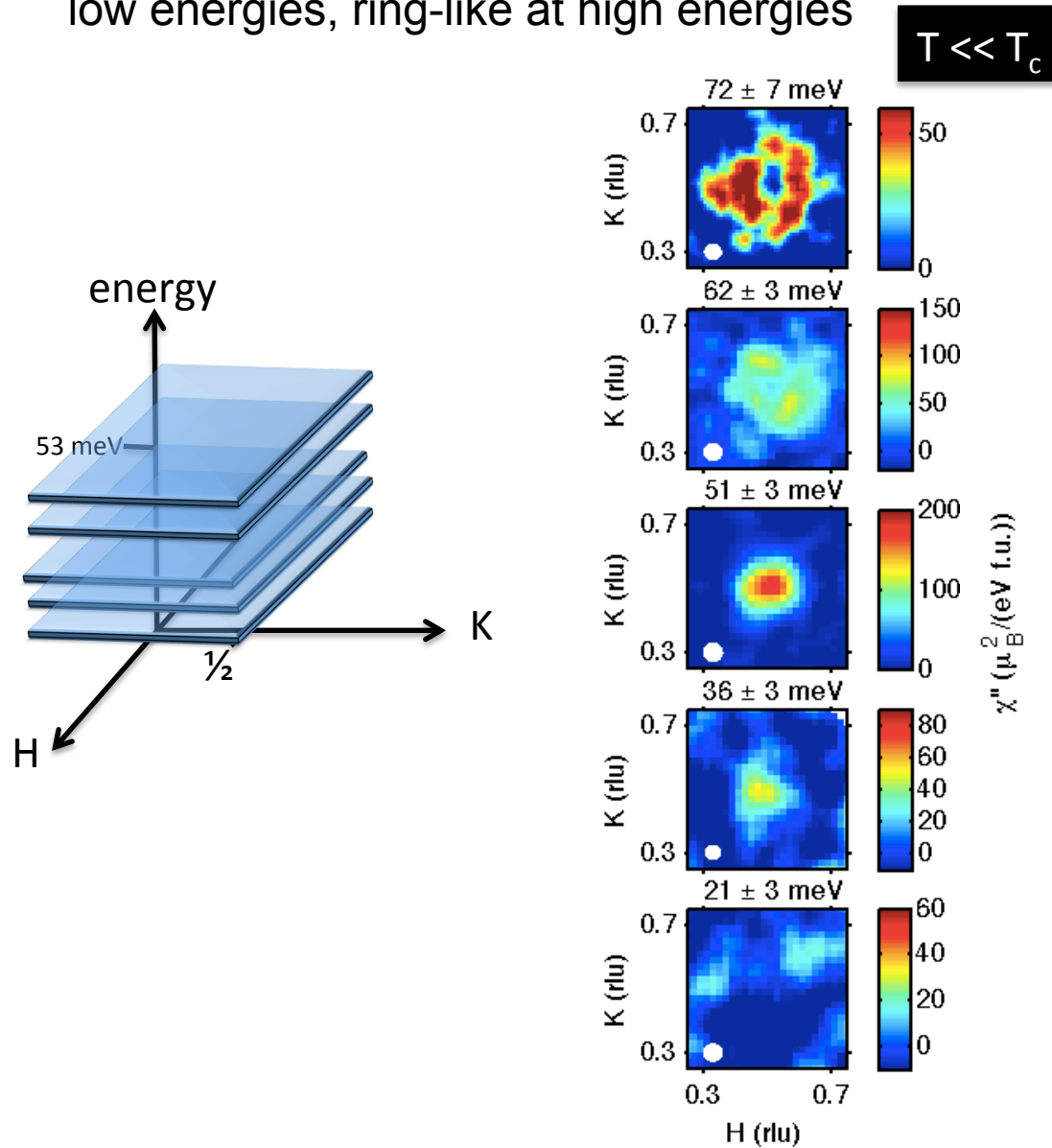


# Antiferromagnetic spectrum



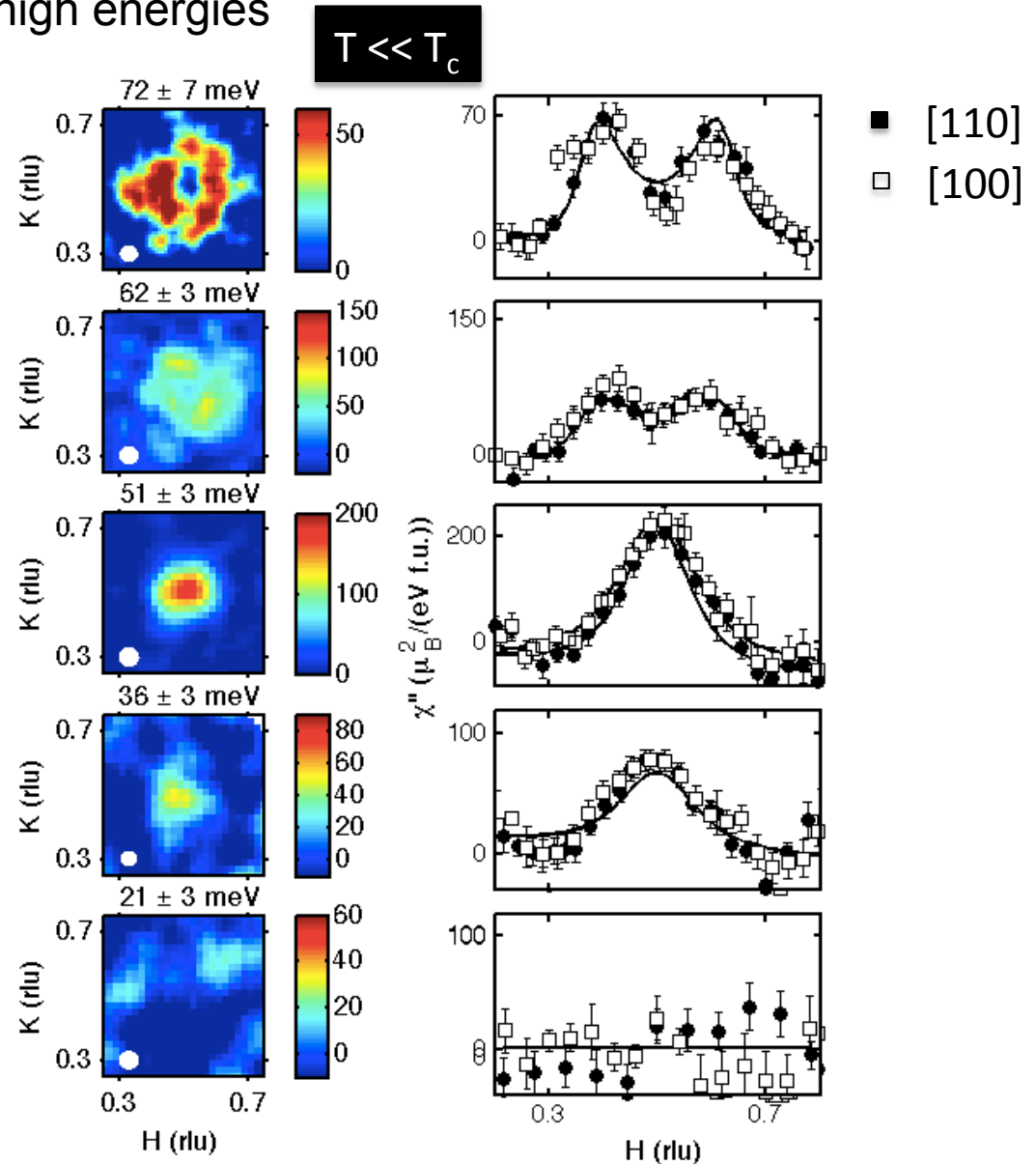
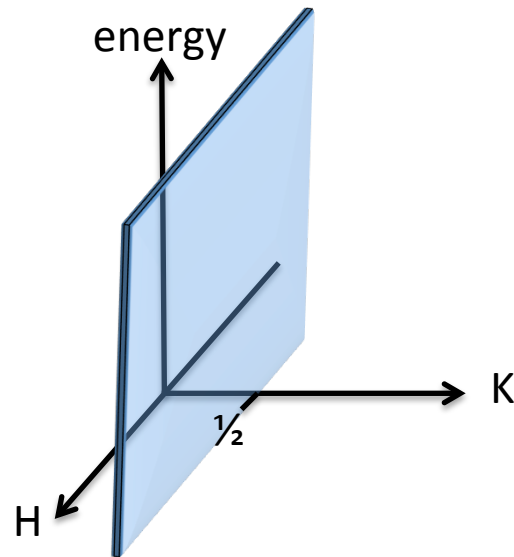
# Antiferromagnetic spectrum

- Constant energy images: reveals the Q-structure: commensurate at low energies, ring-like at high energies



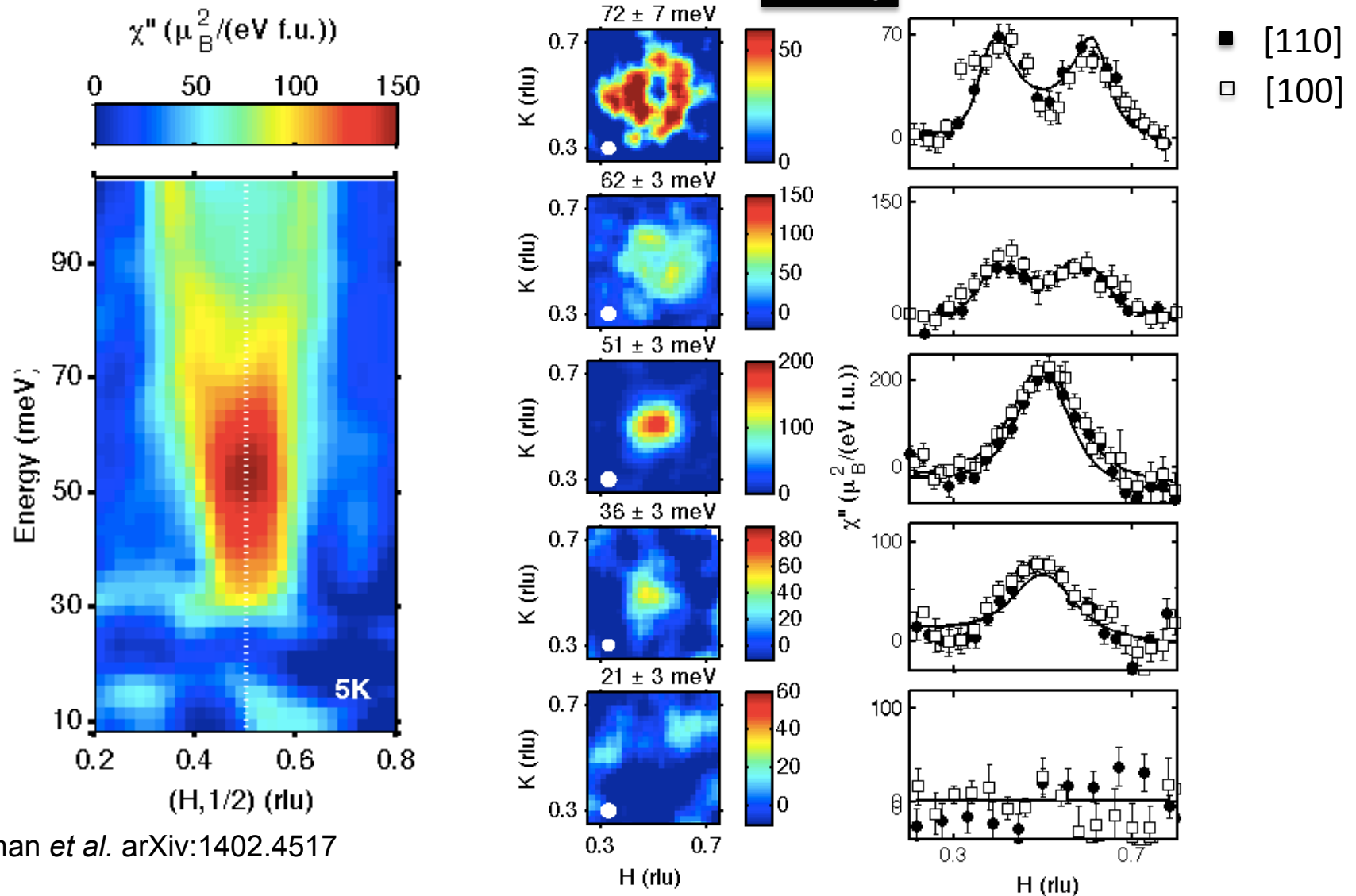
# Antiferromagnetic spectrum

- Constant energy images: reveals the Q-structure: commensurate at low energies, ring-like at high energies



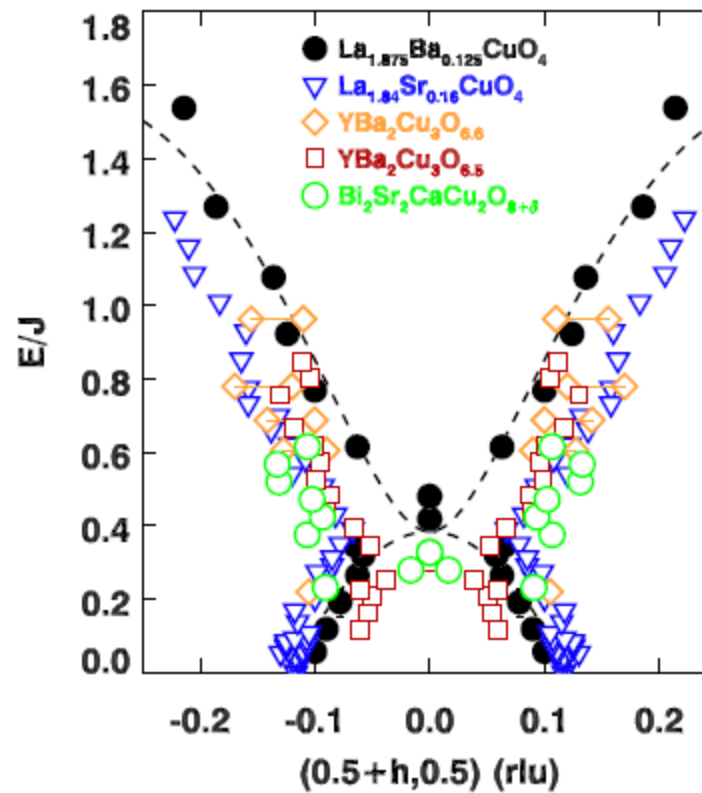
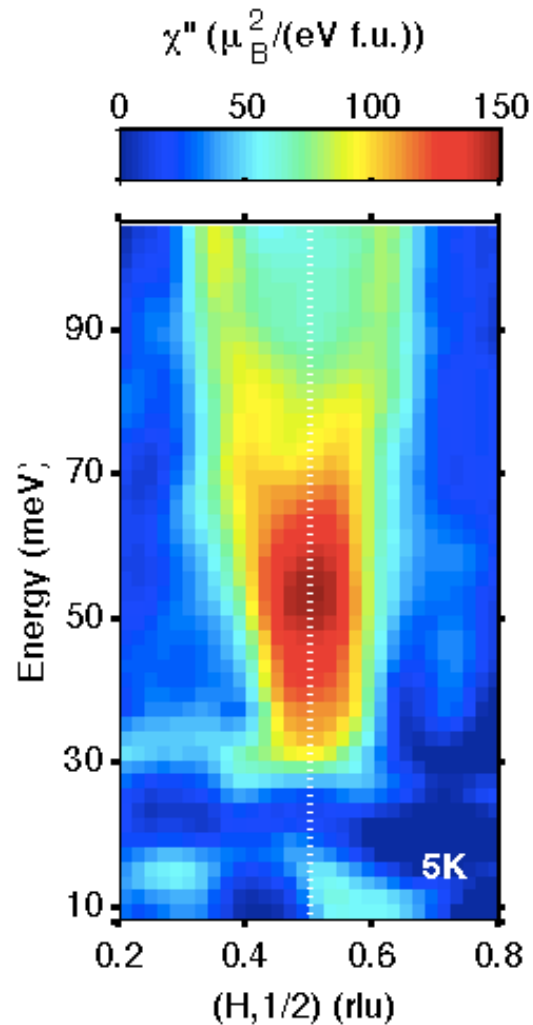
# Antiferromagnetic spectrum

- Constant energy images: reveals the Q-structure: commensurate at low energies, ring-like at high energies



# Antiferromagnetic spectrum

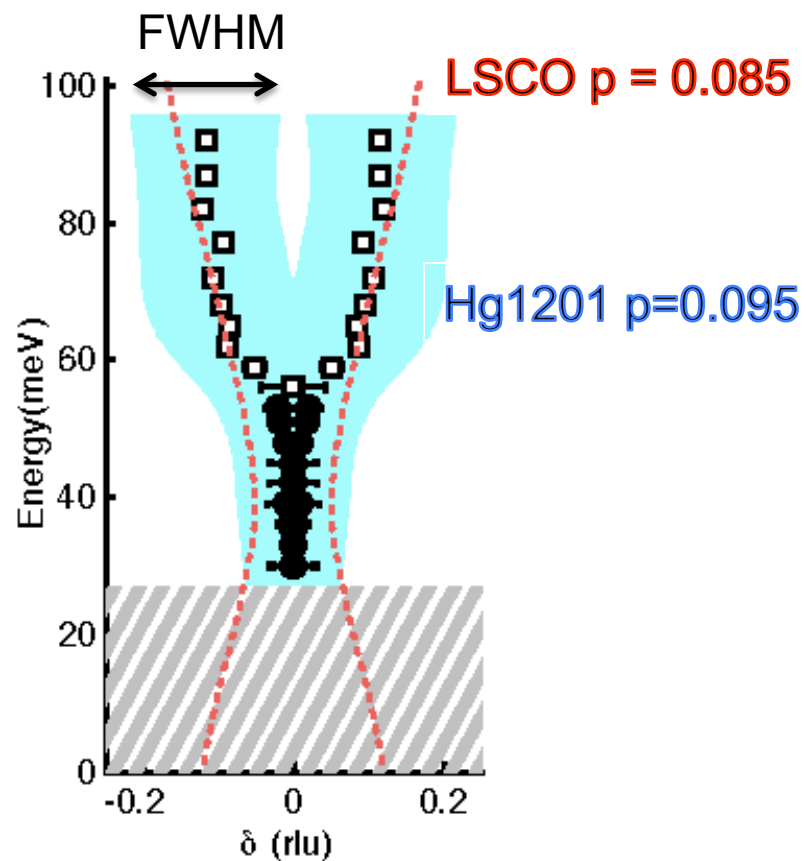
- 'Y'- shaped dispersion with 27 meV gap





## Compare to LSCO

- Gaped vs ungaped
- Low energy commensurate vs incommensurate dispersion
- Stripe correlations, if present would have to be very weak



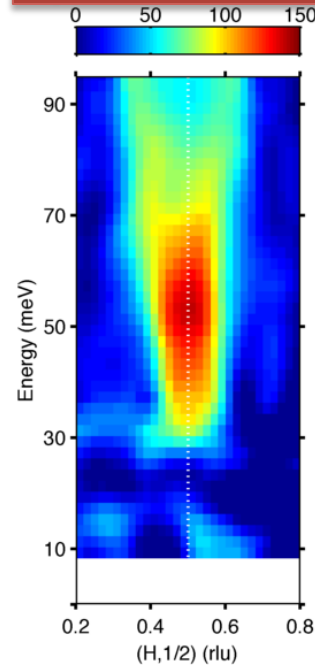
O. J. Lipscombe, *et al.* PRL (2009)

# Doping dependence

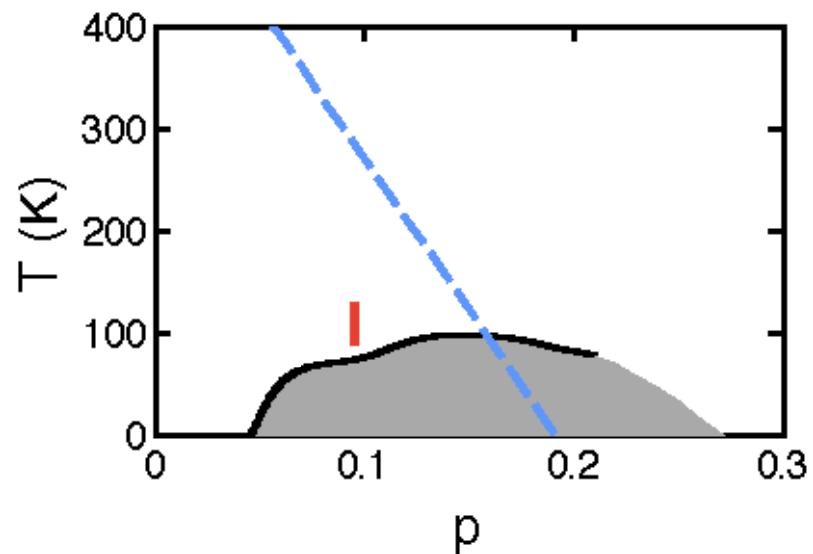
$T_c = 55 \text{ K}; p \approx 0.063$

$T_c = 71 \text{ K}; p \approx 0.095$

$T_c = 88 \text{ K}; p \approx 0.117$



$T = 5 \text{ K}$

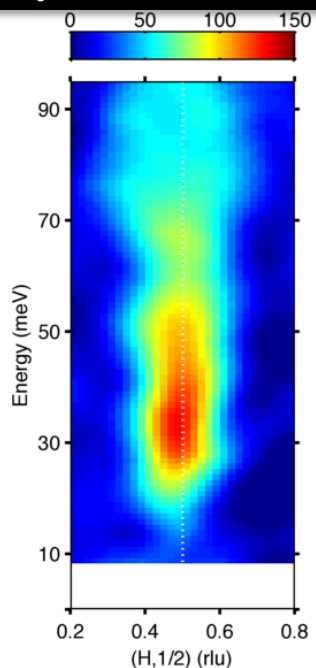


M. K. Chan *et al.* unpublished

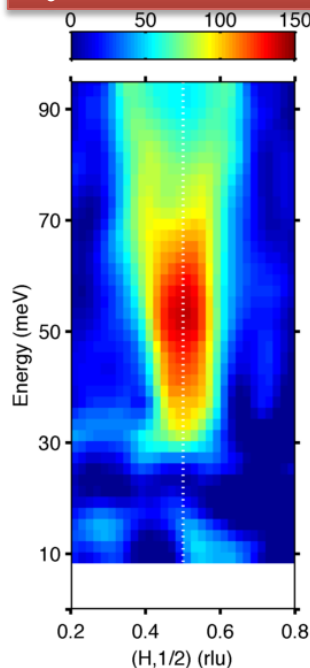
# Doping dependence

$T = 5 \text{ K}$

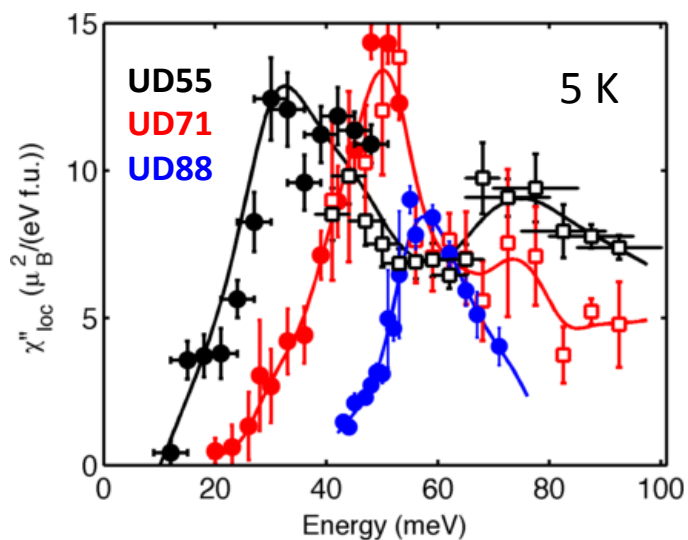
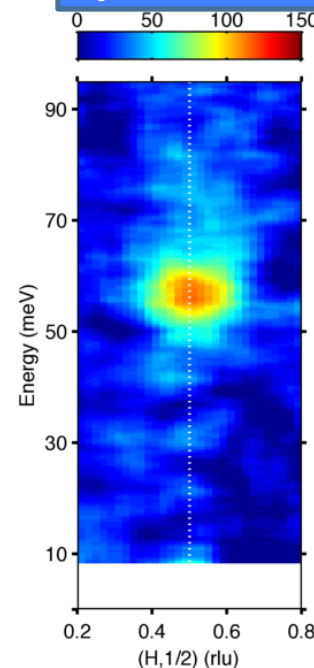
$T_c = 55 \text{ K}; p \approx 0.063$



$T_c = 71 \text{ K}; p \approx 0.095$



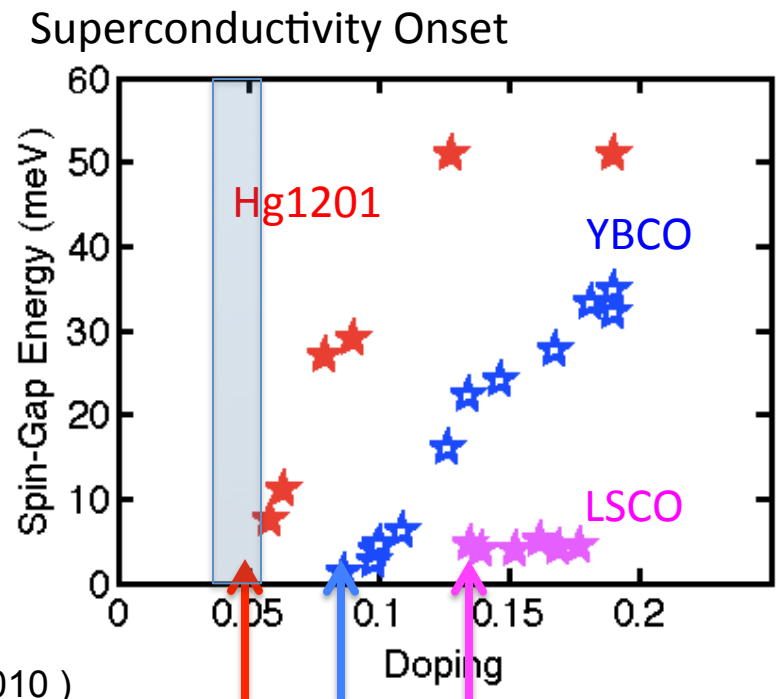
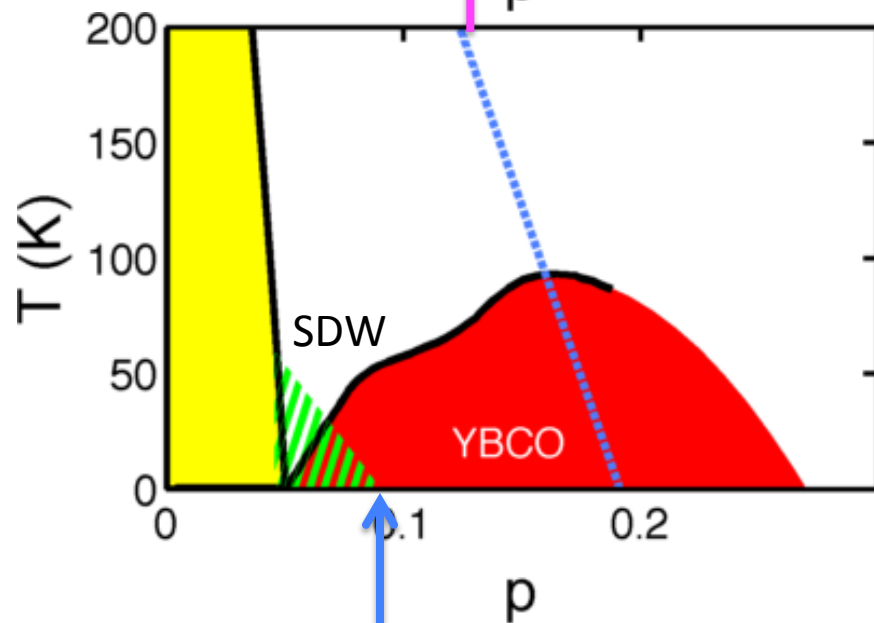
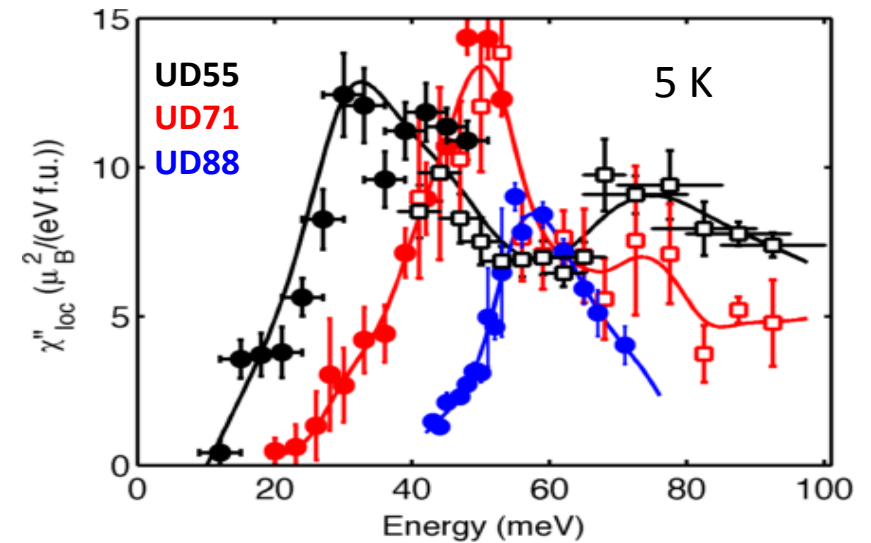
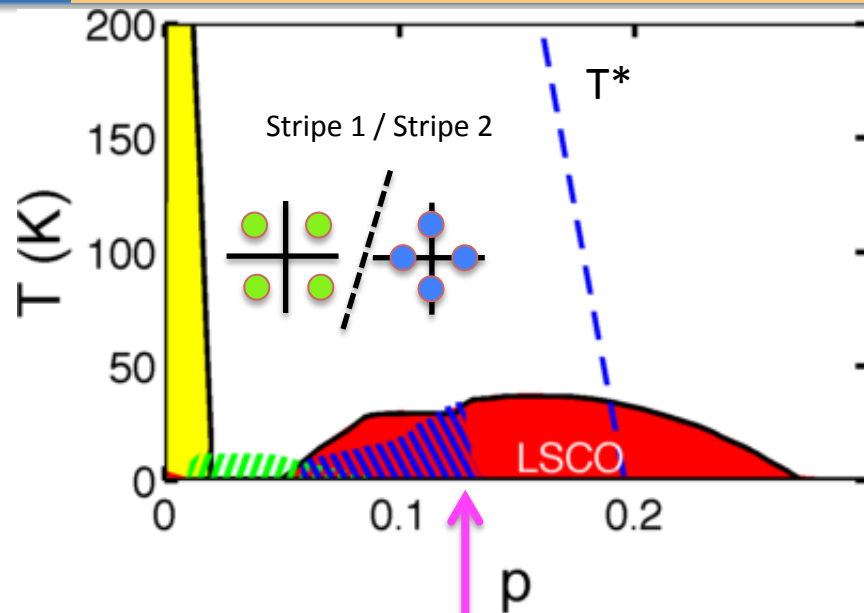
$T_c = 88 \text{ K}; p \approx 0.117$



- Magnetic weight decreases with increasing doping

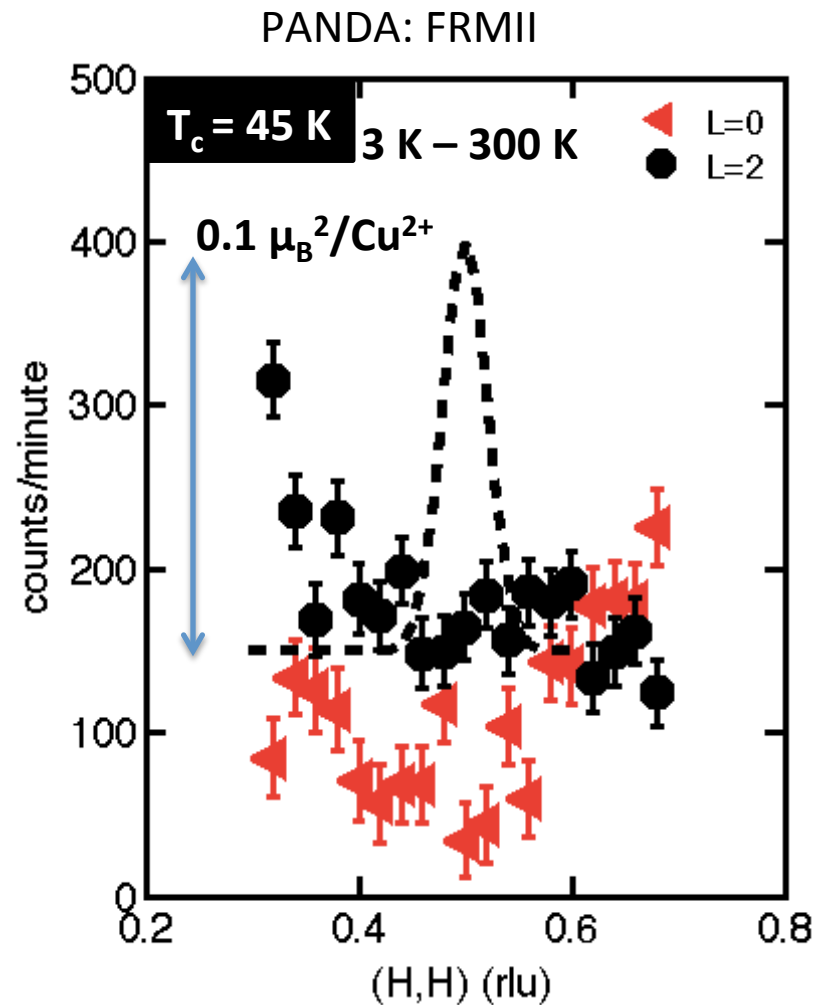
M. K. Chan *et al.* unpublished

# What is the magnetic phase diagram?



P. Bourges Physica B (1995);  
 M. Kofu *et al.* PRL (1999); D. Haug *et al.* New Jour. of Phys. (2010 )

## No static magnetism



- No commensurate or diagonal elastic response of the same size as in LSCO or YBCO

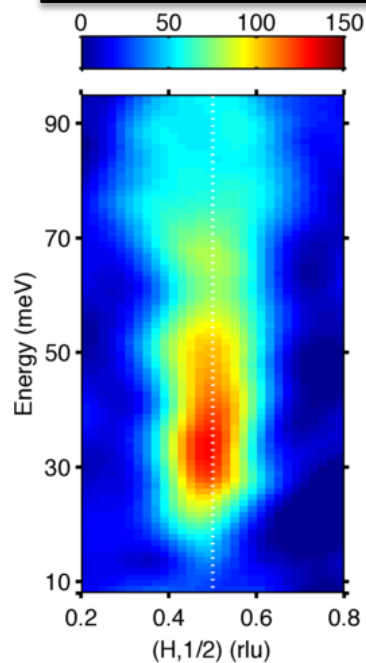
## Neutron: Conclusion 1

- Low energy **commensurate** fluctuations: weak or absent stripe correlations
- No static local magnetism in the most underdoped Hg1201 sample studied ( $p = 0.058$ )

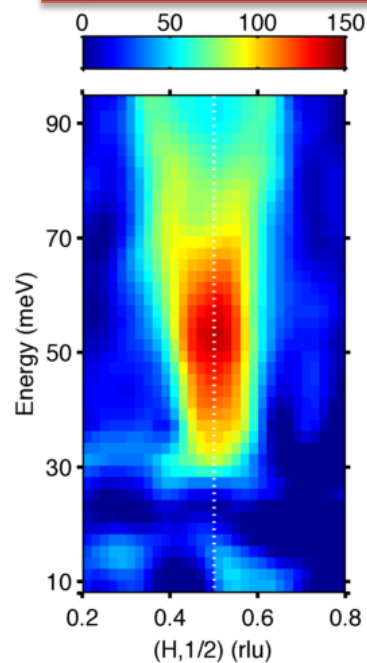
# Change across $T^*$ (pseudogap)

$T = 5 \text{ K}$

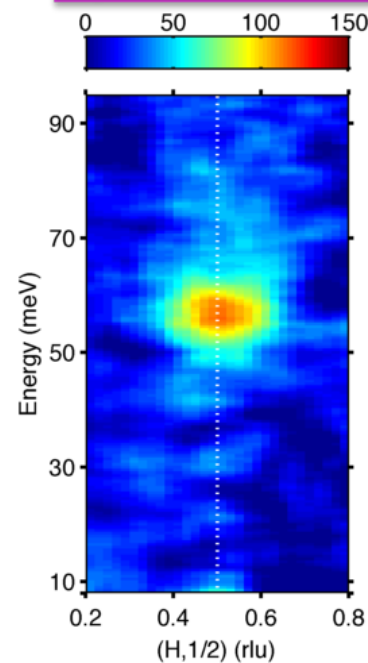
$T_c = 55 \text{ K}; p \approx 0.063$



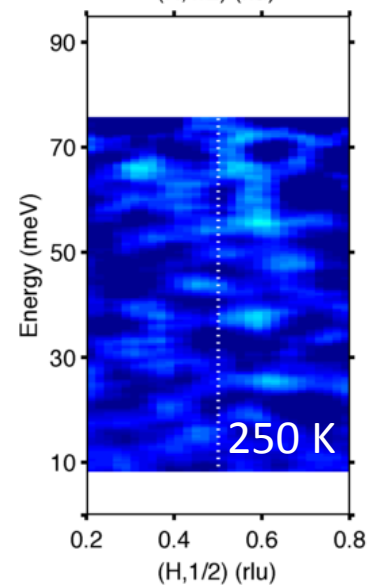
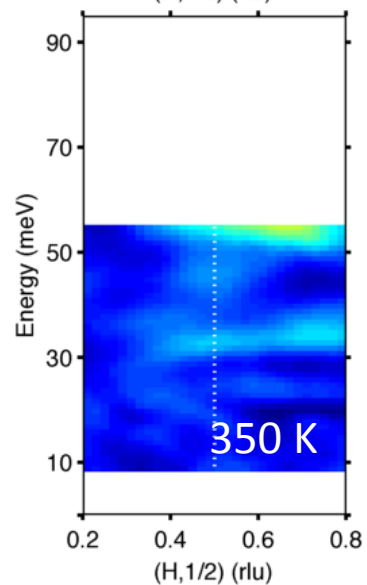
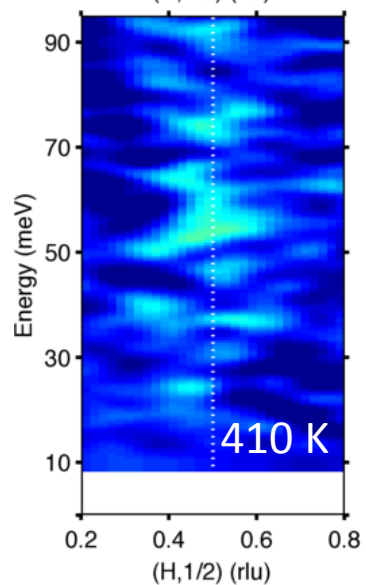
$T_c = 71 \text{ K}; p \approx 0.095$



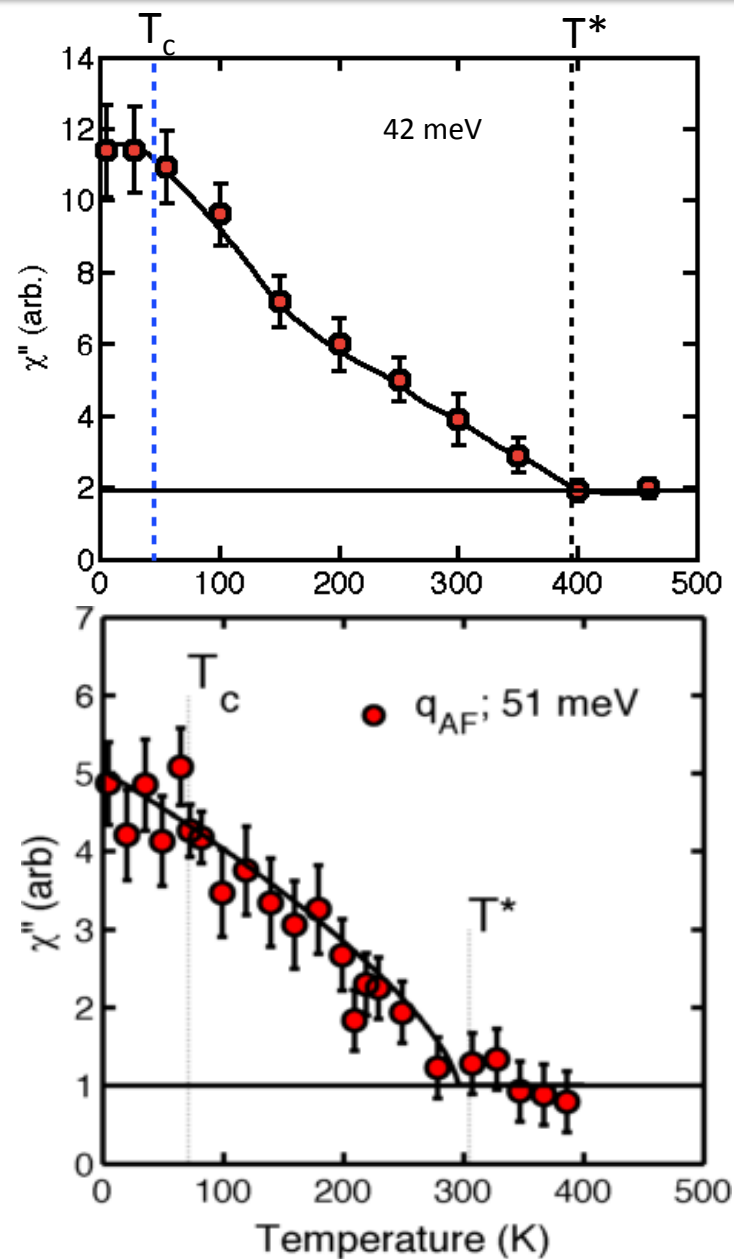
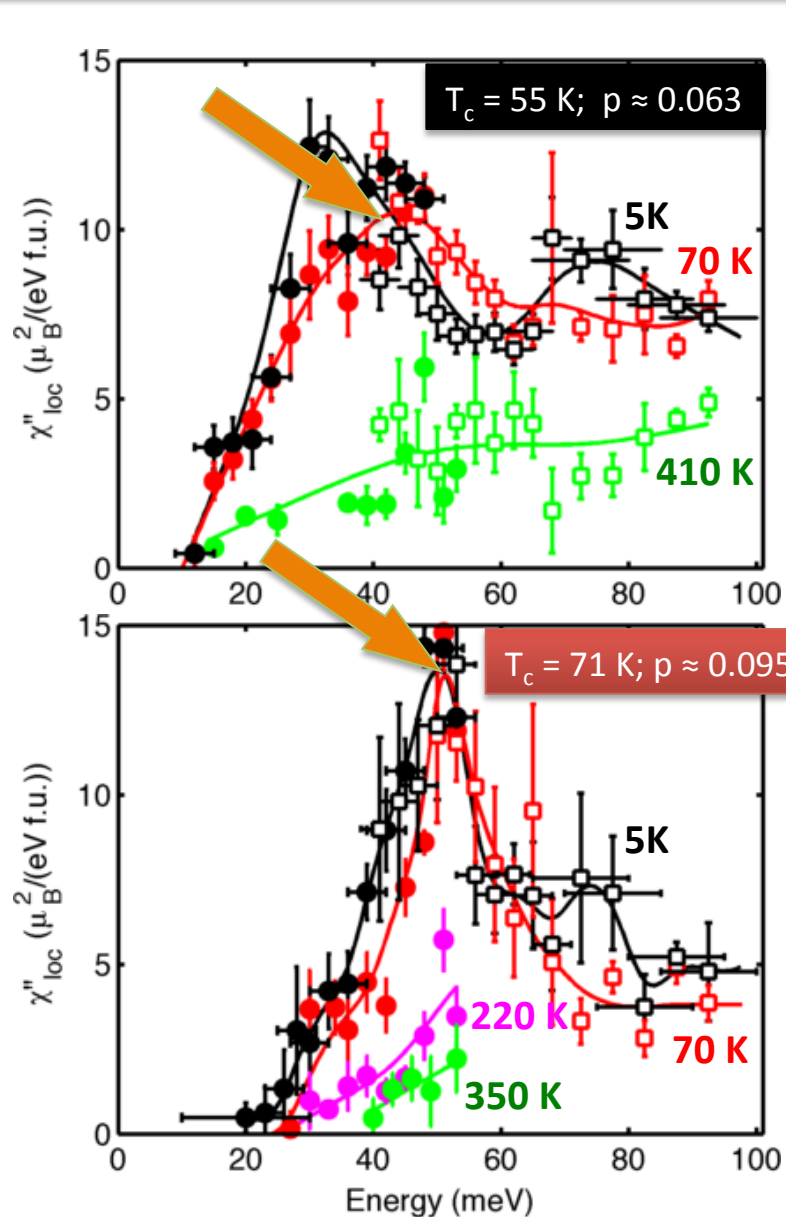
$T_c = 88 \text{ K}; p \approx 0.117$



$T > T^*$

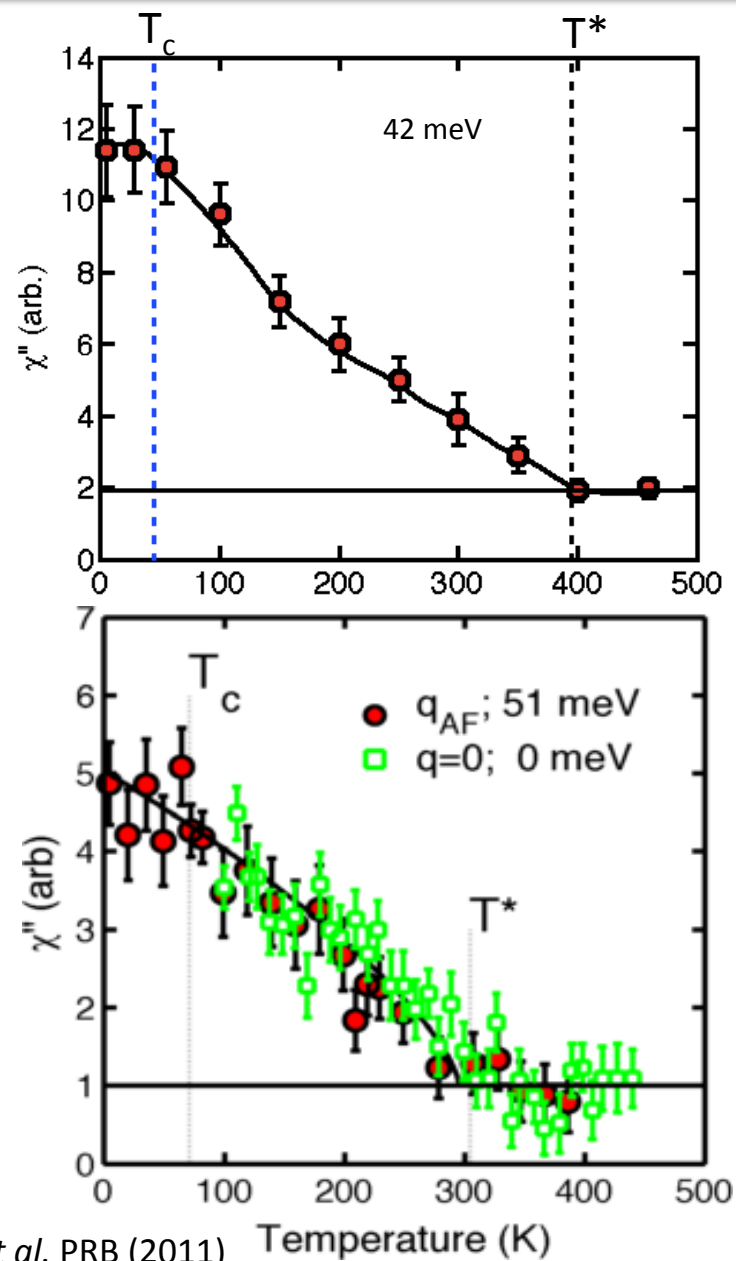
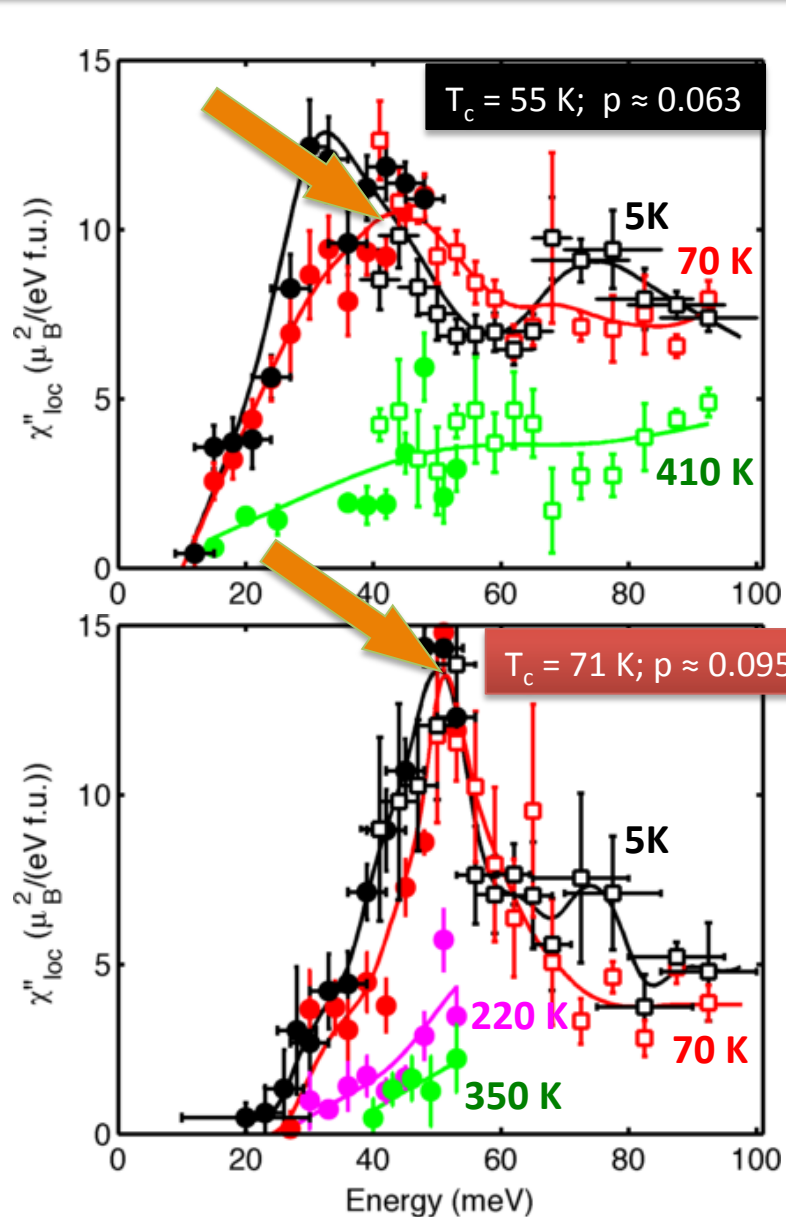


# Onset of AF fluctuations





# Onset of AF fluctuations



Y. Li. *et al.* PRB (2011)

## Neutron: Conclusion 2

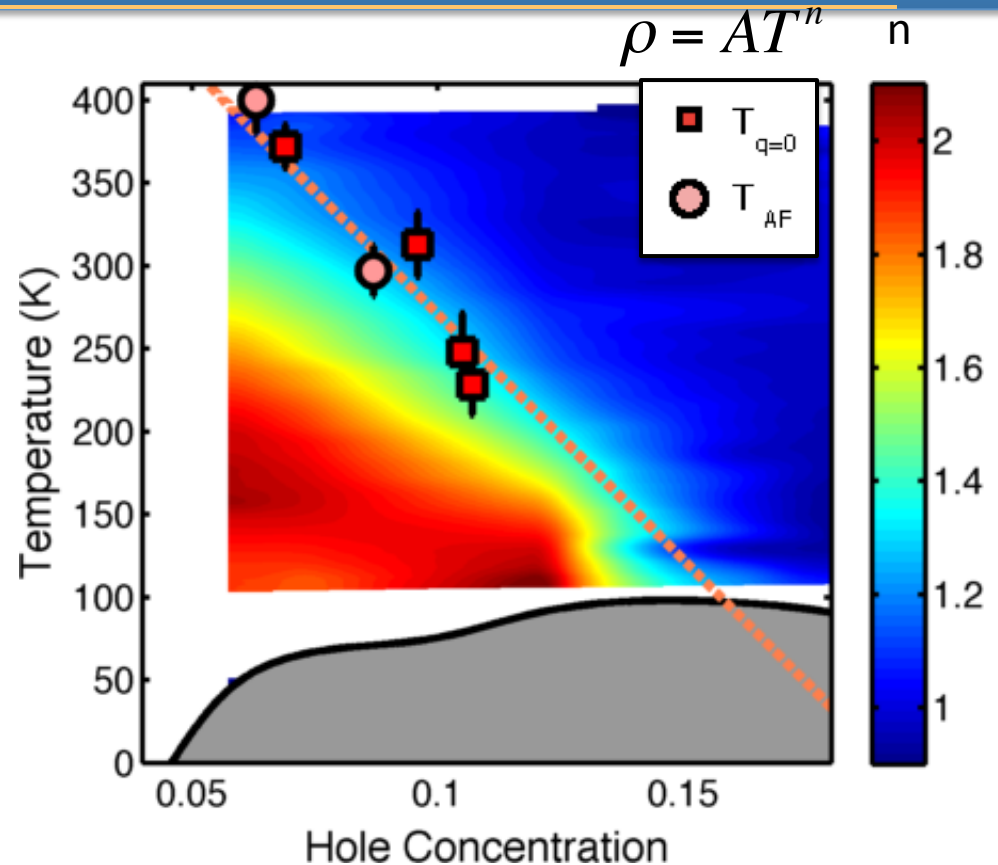
- Low energy commensurate fluctuations: weak or absent stripe correlations
- No static local magnetism in the most underdoped Hg1201 sample studied ( $p = 0.058$ ) : supported by resonance and a spin gap
- Majority of fluctuation weight onsets at the pseudogap temperature  $T^*$

# Conclusions

Hg1201: highest  $T_c$ , lower disorder,  
simple tetragonal structure

Pseudogap characterized by:

- Fermi-liquid like electrical transport with a largely stable Fermi-surface (w.r.t. temperature)
- Onset of **commensurate** magnetic fluctuations
- No incommensurate AF ordering

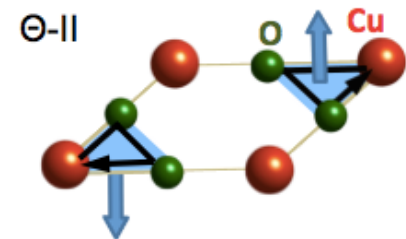


B. Fauqué *et al.* PRL (2006)

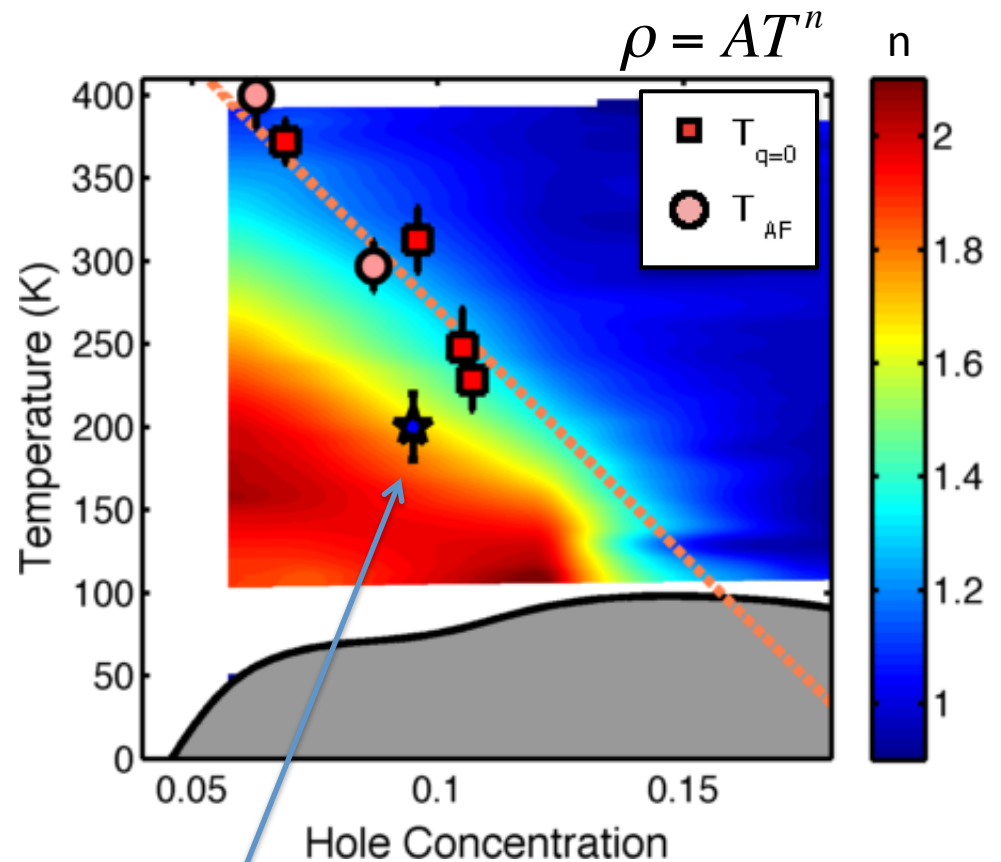
Y. Li. *et al.* PRB (2011)

N. Barišić, MKC *et al.* PNAS (2013)

$q=0$  loop current order



# Cuprate phase diagram controlled by AF fluctuations



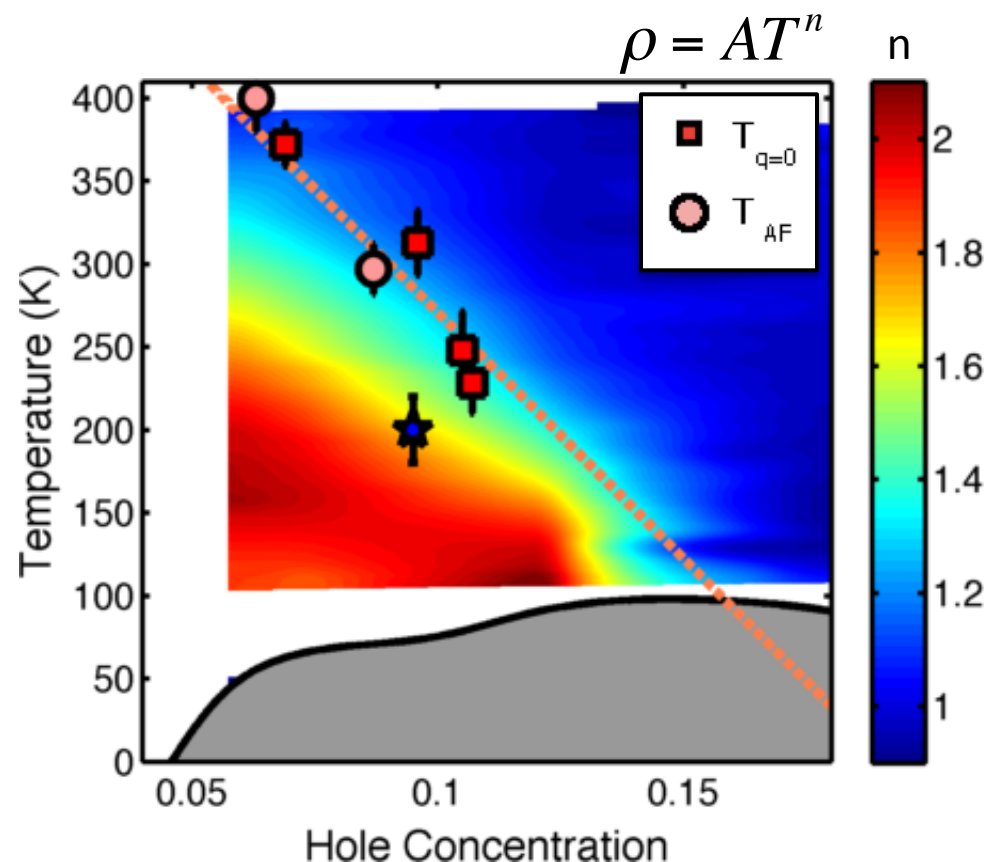
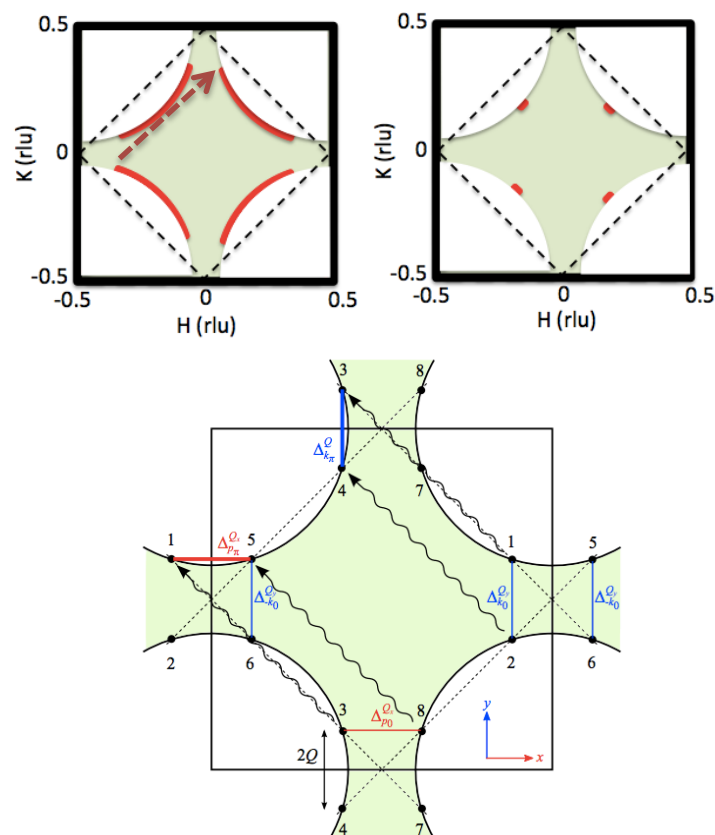
W. Tabis *et al.* arXiv:1404.7658

Short-ranged CDW observed in Hg1201

Ghiringhelli *et al.*, Science (2012)

Chang *et al.*, Nature Phys. (2012)

# Cuprate phase diagram controlled by AF fluctuations



Y. Wang & A. Chubukov arxiv:1401.0712  
M. A. Metlitski & S. Sachdev, PRB (2010).

- Recent and old theoretical proposals based on spin fluctuations
- the hot spots (connected by  $q_{AF}$ ) near the anti-nodes
- Nodal quasiparticles are Fermi-liquid like